

Exotic nuclei with charm and bottom flavors

Shigehiro Yasui
KEK

27-28 Nov. 2009

新学術領域

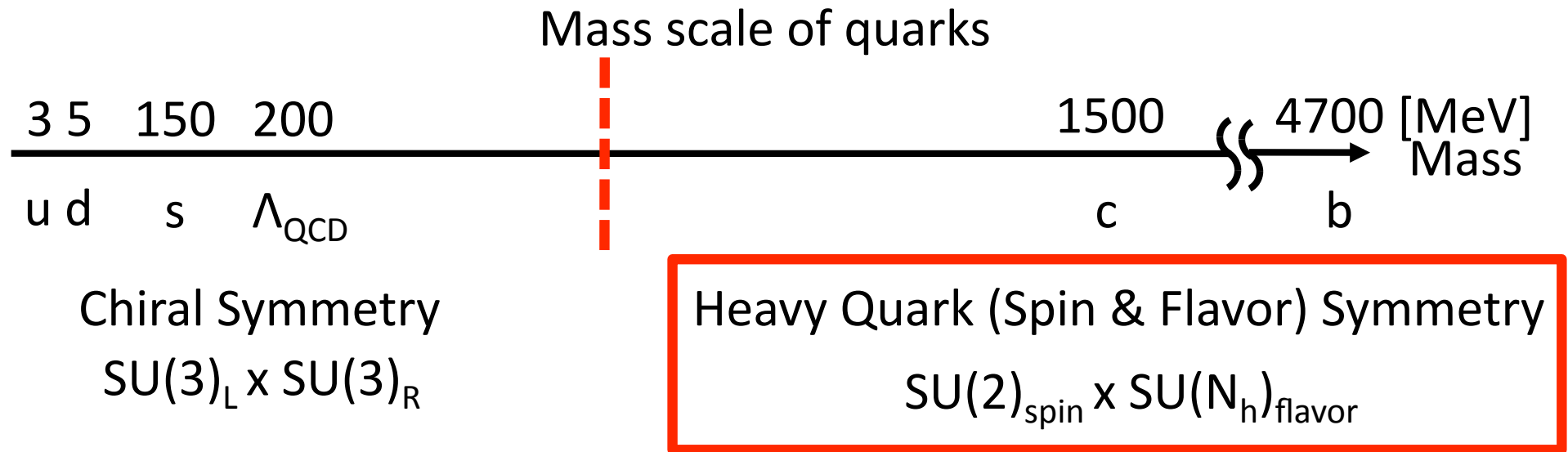
「多彩なフレーバーで探る新しいハドロン存在形態の包括的研究」

キックオフミーティング@Nagoya University

Contents

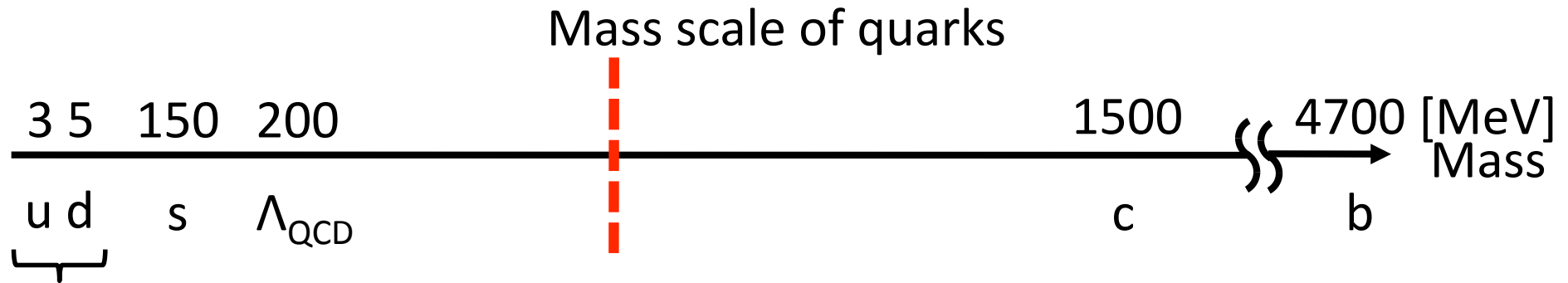
1. Motivation
2. DN and BN potential
 - Heavy Quark Symmetry
3. $D^{\text{bar}}N$ and BN bound states
4. Exotic charm/bottom nuclei
 - $D^{\text{bar}}NN$ and BNN
5. Conclusion & Perspective

Why are heavy flavors interesting in QCD?



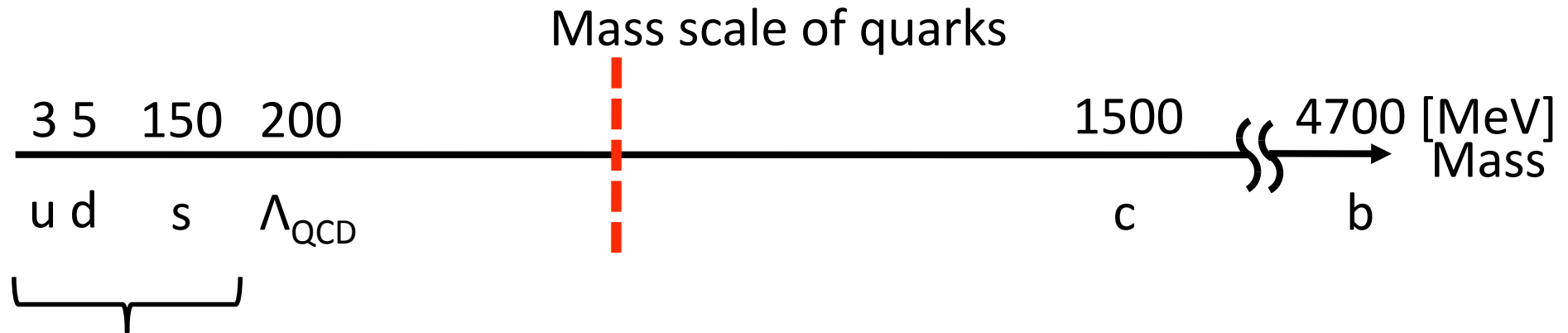
- ✓ New energy scale ($m_{c,b}$) is introduced in addition to Λ_{QCD} .
 - ✓ New symmetry (heavy quark symmetry) appears.
 - ✓ QCD Effective Theory (HQET, NRQCD, pNRQCD, ...)
 - ✓ Lattice QCD approach
- etc.

Why are heavy flavors interesting in QCD?



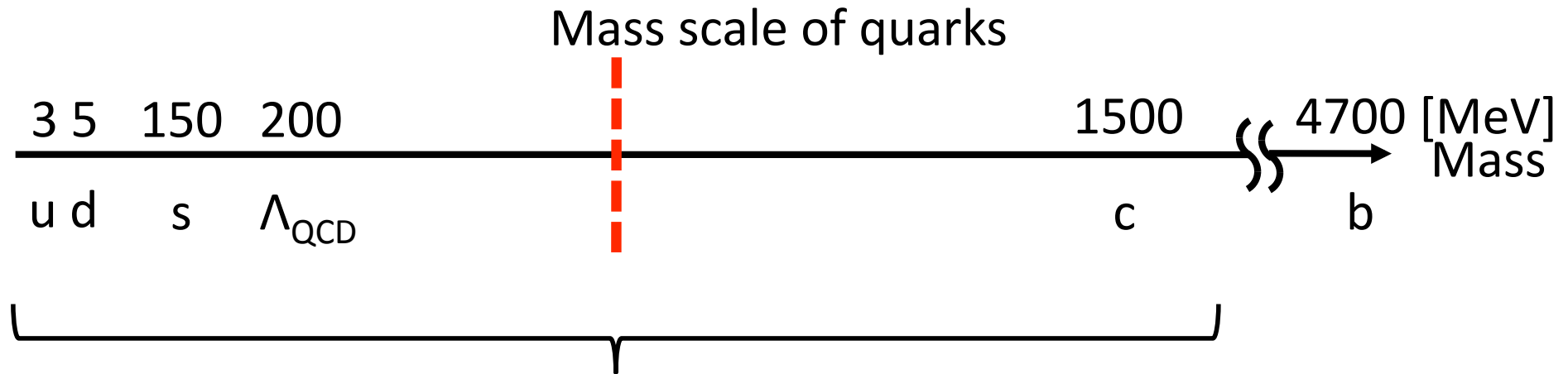
- Normal Nuclei
- Large Isospin Nuclei
(neutron/proton rich)

Why are heavy flavors interesting in QCD?



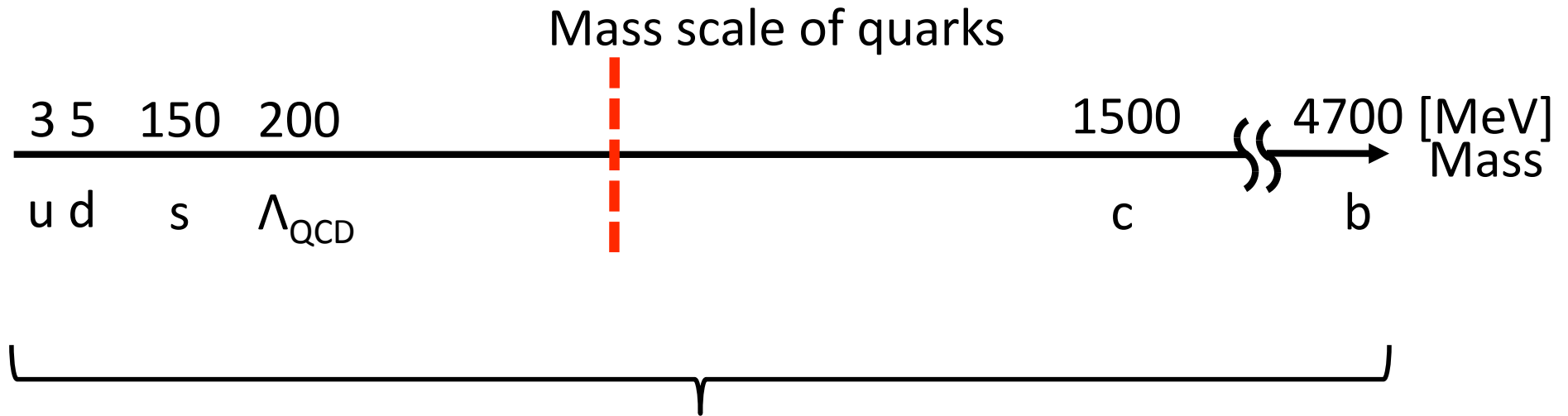
- Strangeness Nuclei @J-PARC
(Hypernuclei/K-nuclei)

Why are heavy flavors interesting in QCD?



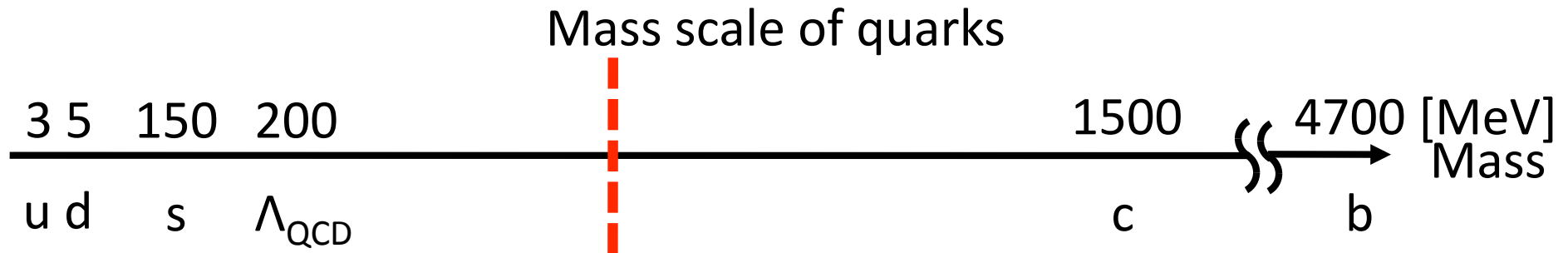
- Charmed Nuclei ? @J-PARC and GSI
(Charm baryon nuclei/D-nuclei)

Why are heavy flavors interesting in QCD?



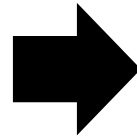
- Bottom Nuclei ??
(Bottom baryon nuclei/B-nuclei)

Why are heavy flavors interesting in QCD?



Theory based on QCD

- Heavy Quark Symmetry
- QCD Effective Theory (HQET, NRQCD, pNRQCD, ...)
- Lattice QCD approach

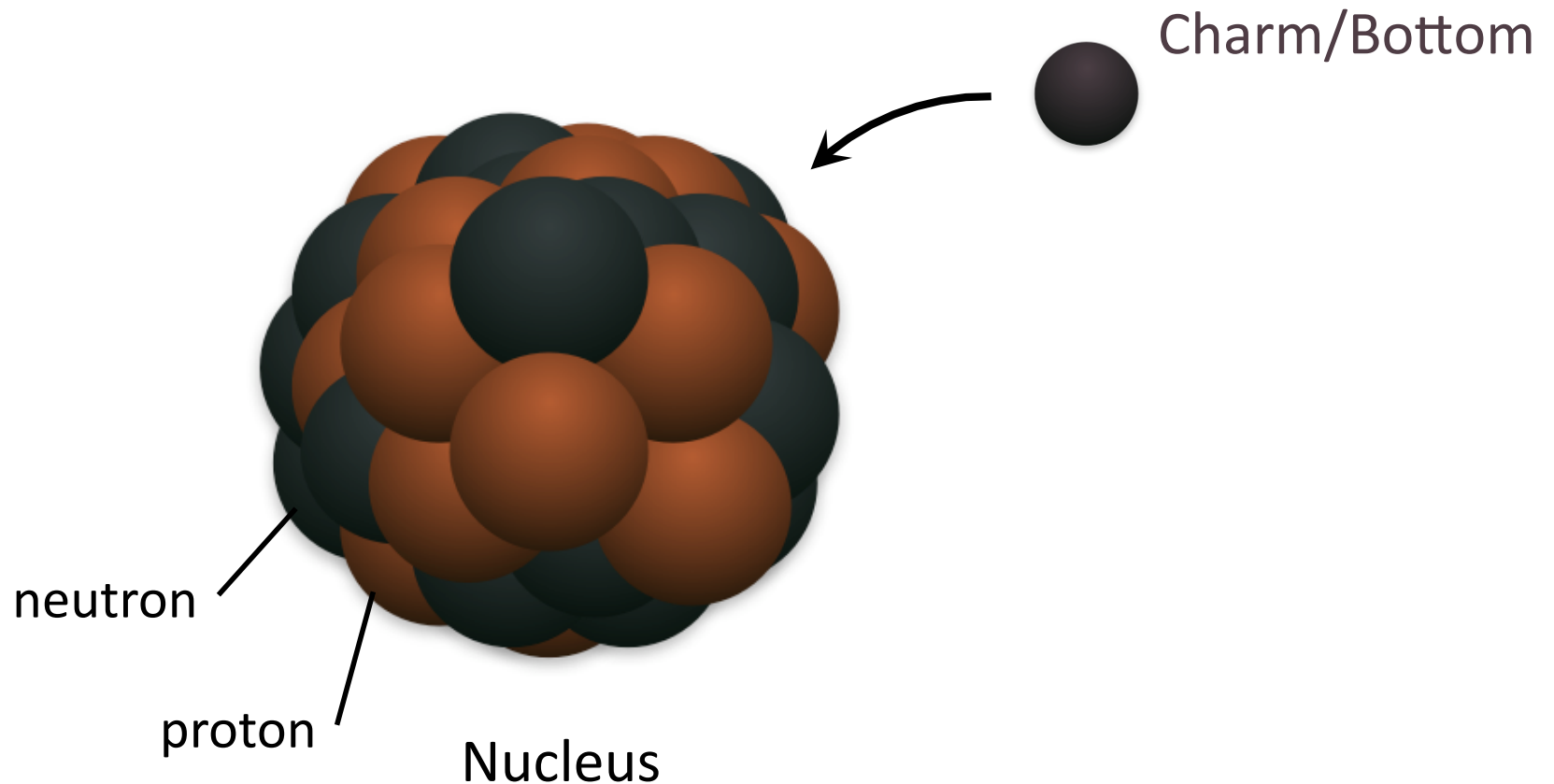


Charm/Bottom Nuclei

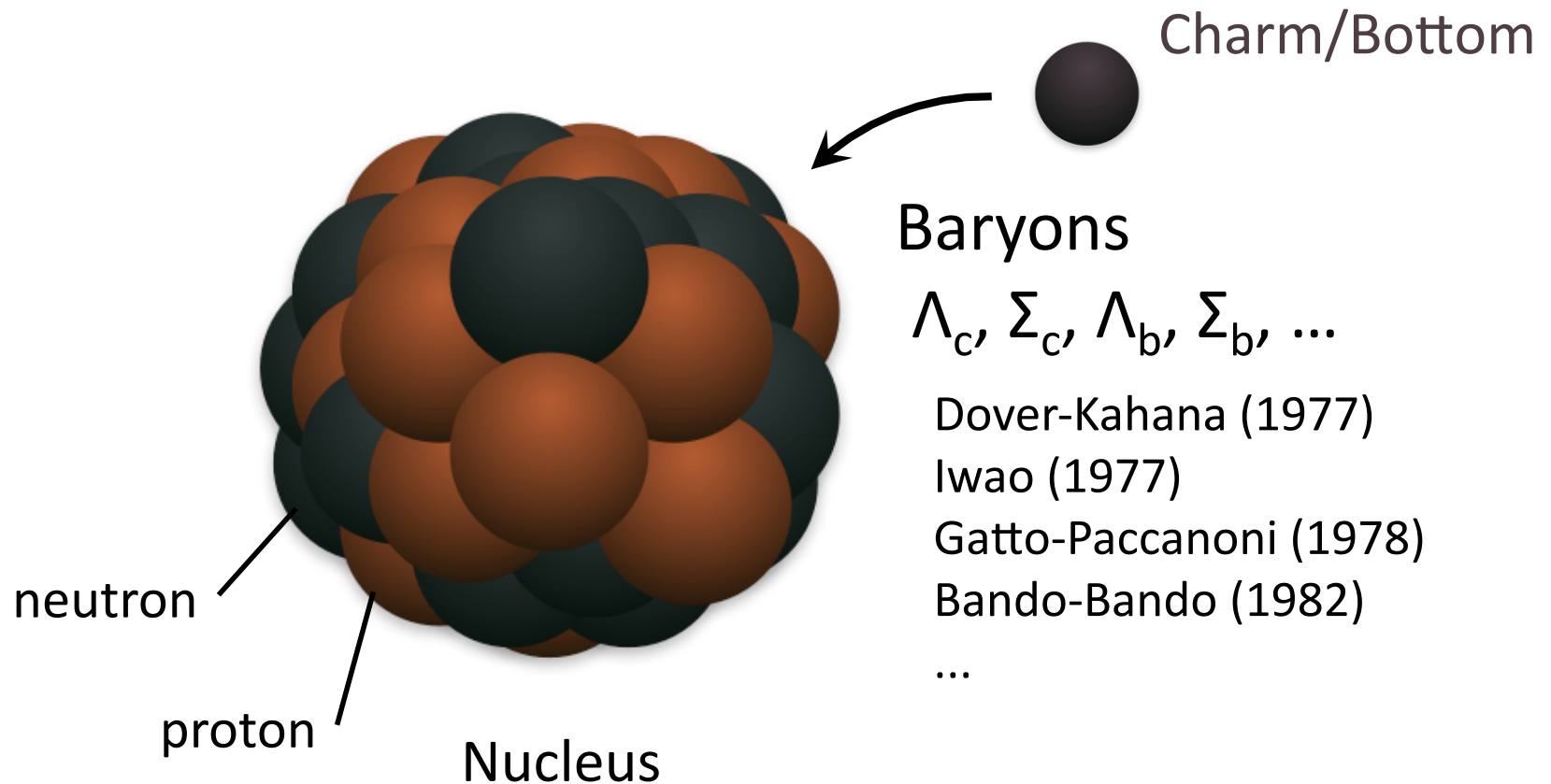
- New states and properties
- Further understanding of exotic nuclei with flavors
- Confinement, χ SB, ...

Cf. Charmonium-Nucleon interaction from Lattice QCD; Hatsuda-Kawanai-Sasaki

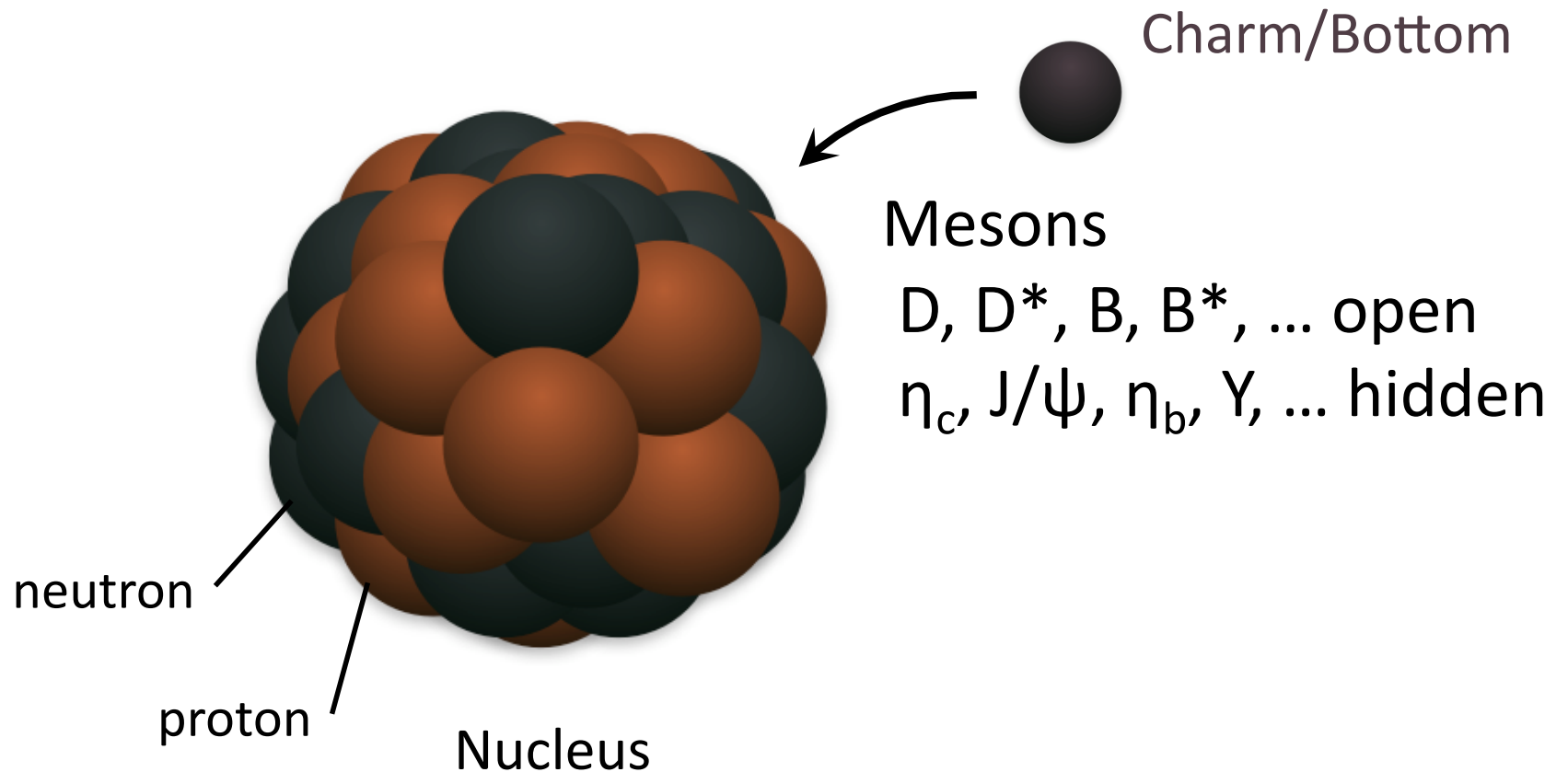
What are charm/bottom nuclei?



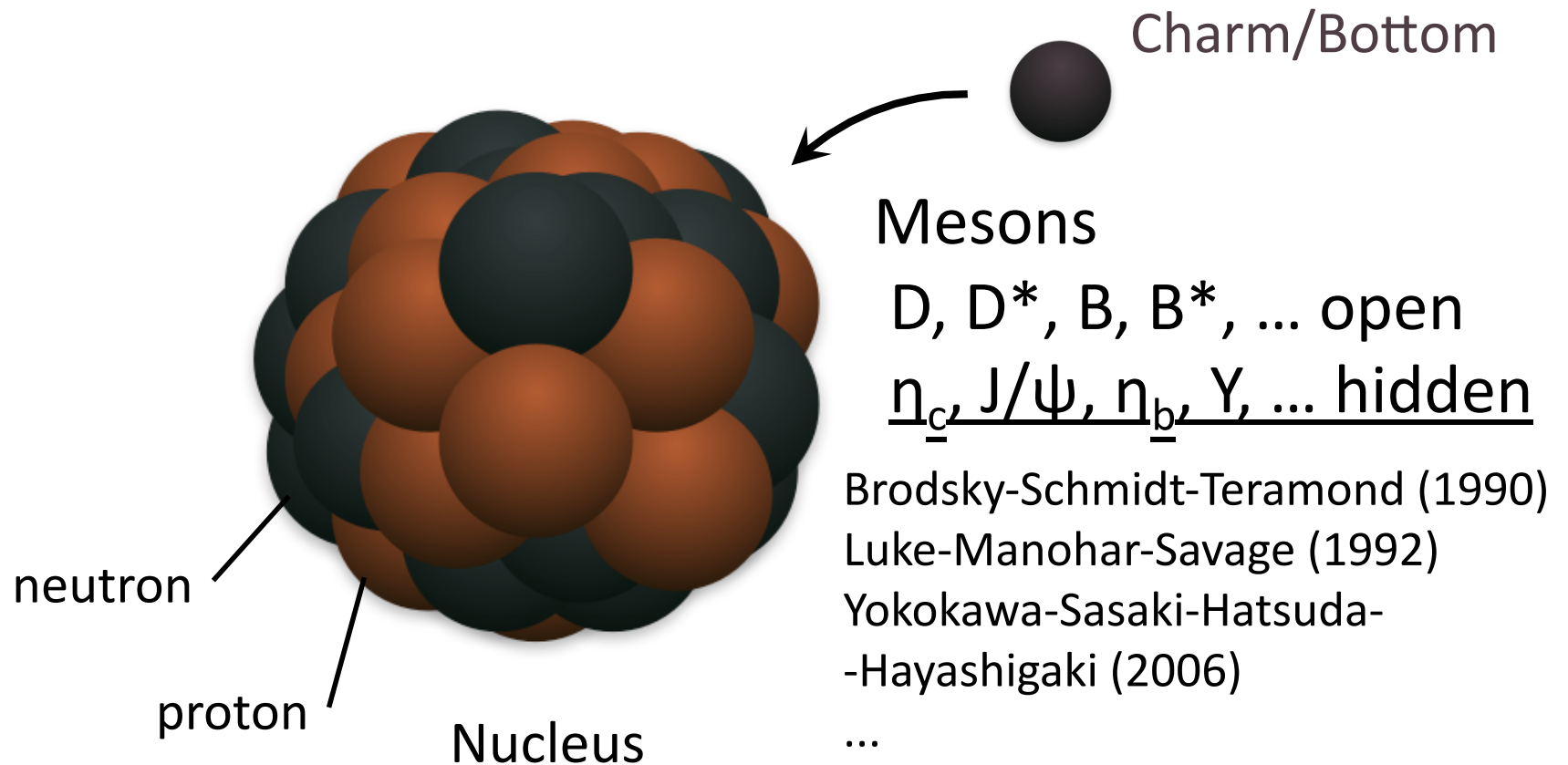
What are charm/bottom nuclei?



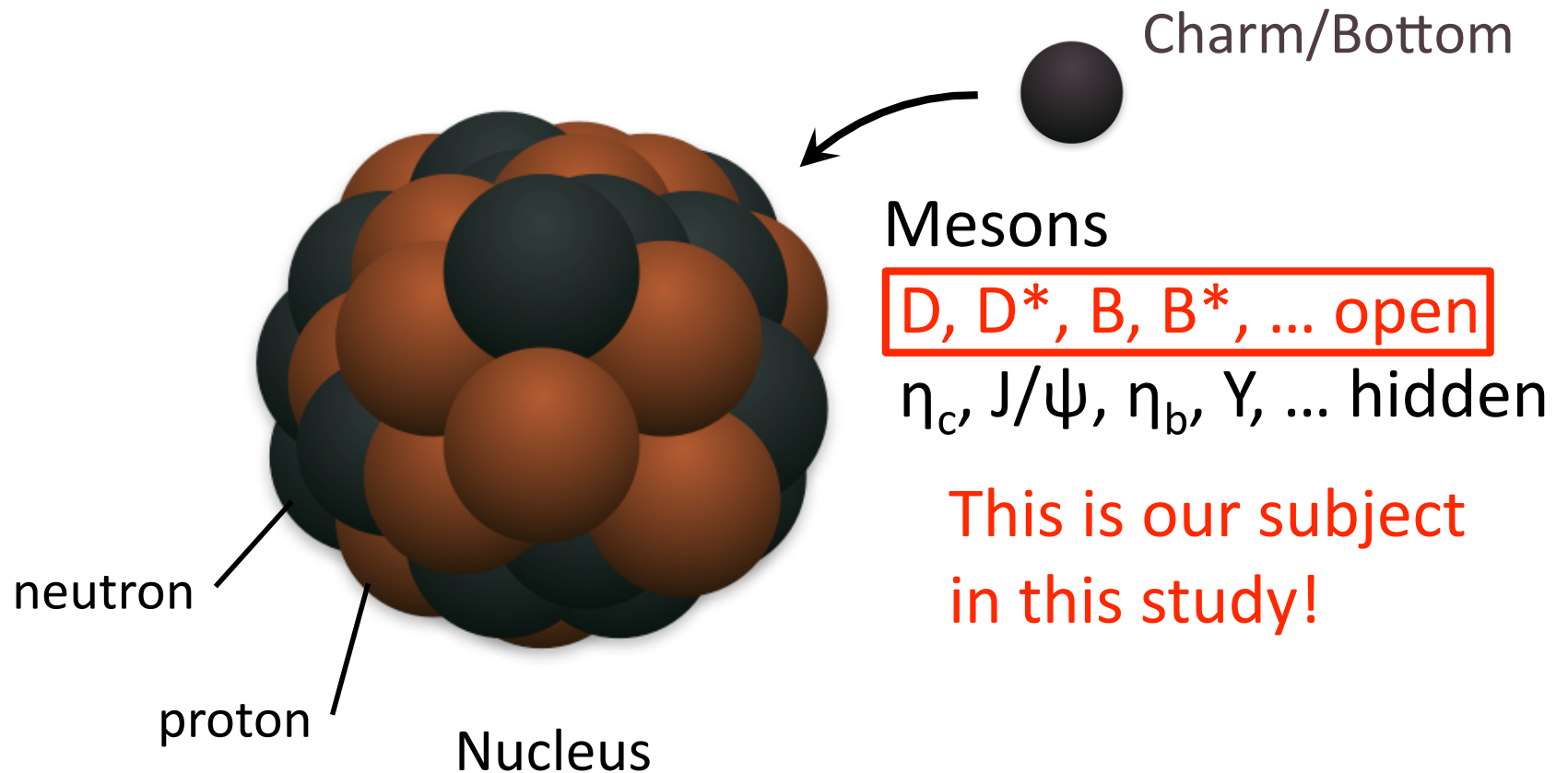
What are charm/bottom nuclei?



What are charm/bottom nuclei?

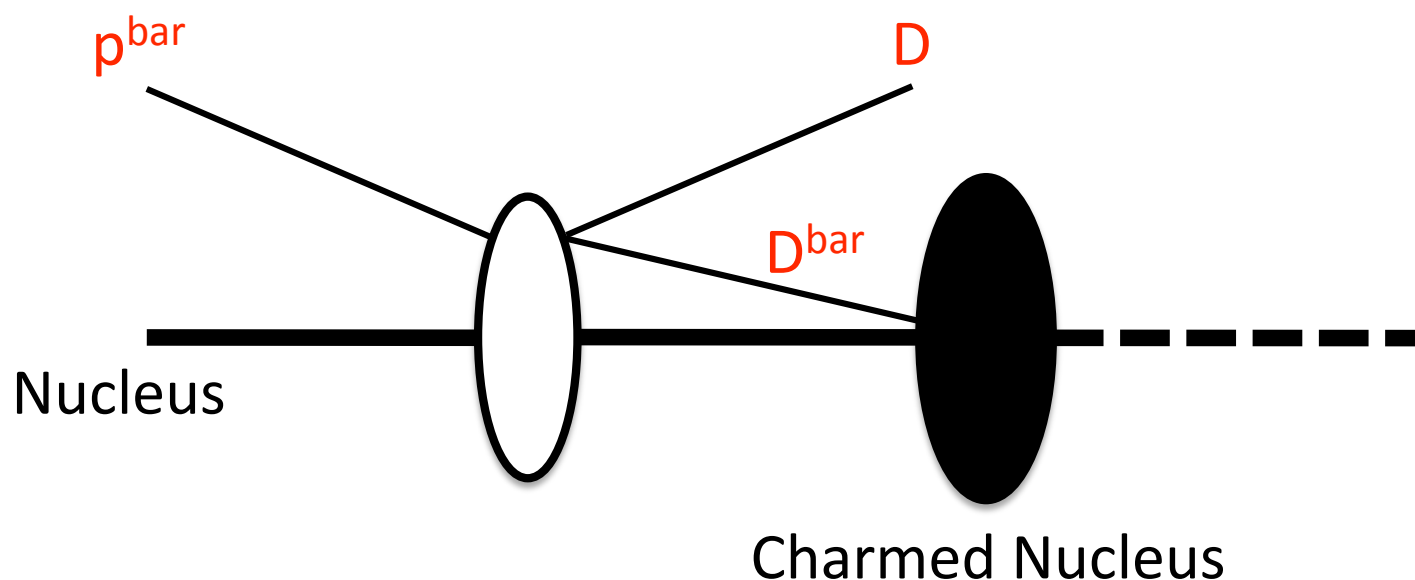


What are charm/bottom nuclei?



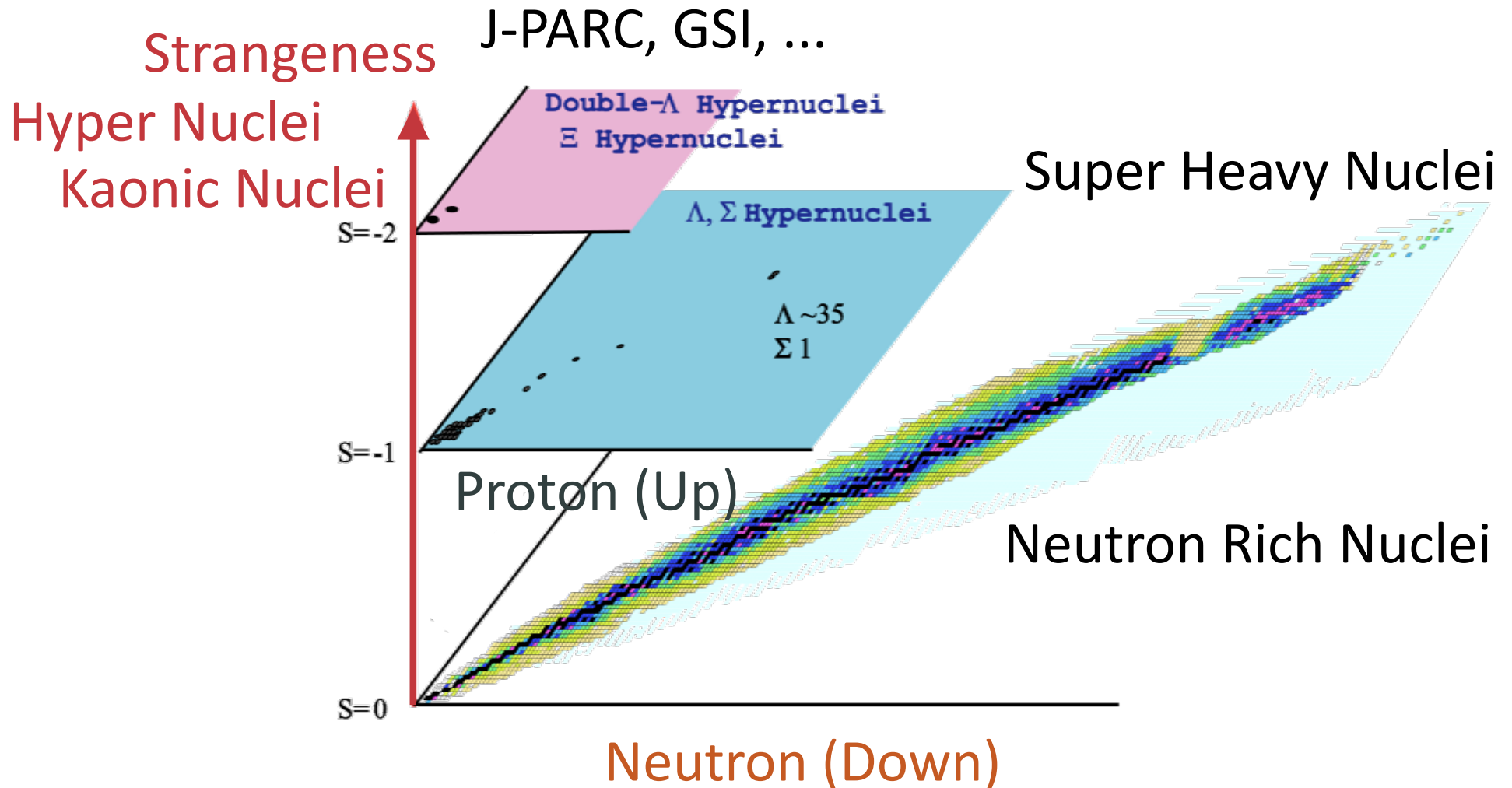
How to generate charm nuclei?

Possible experiments in J-PARC, GSI, ...

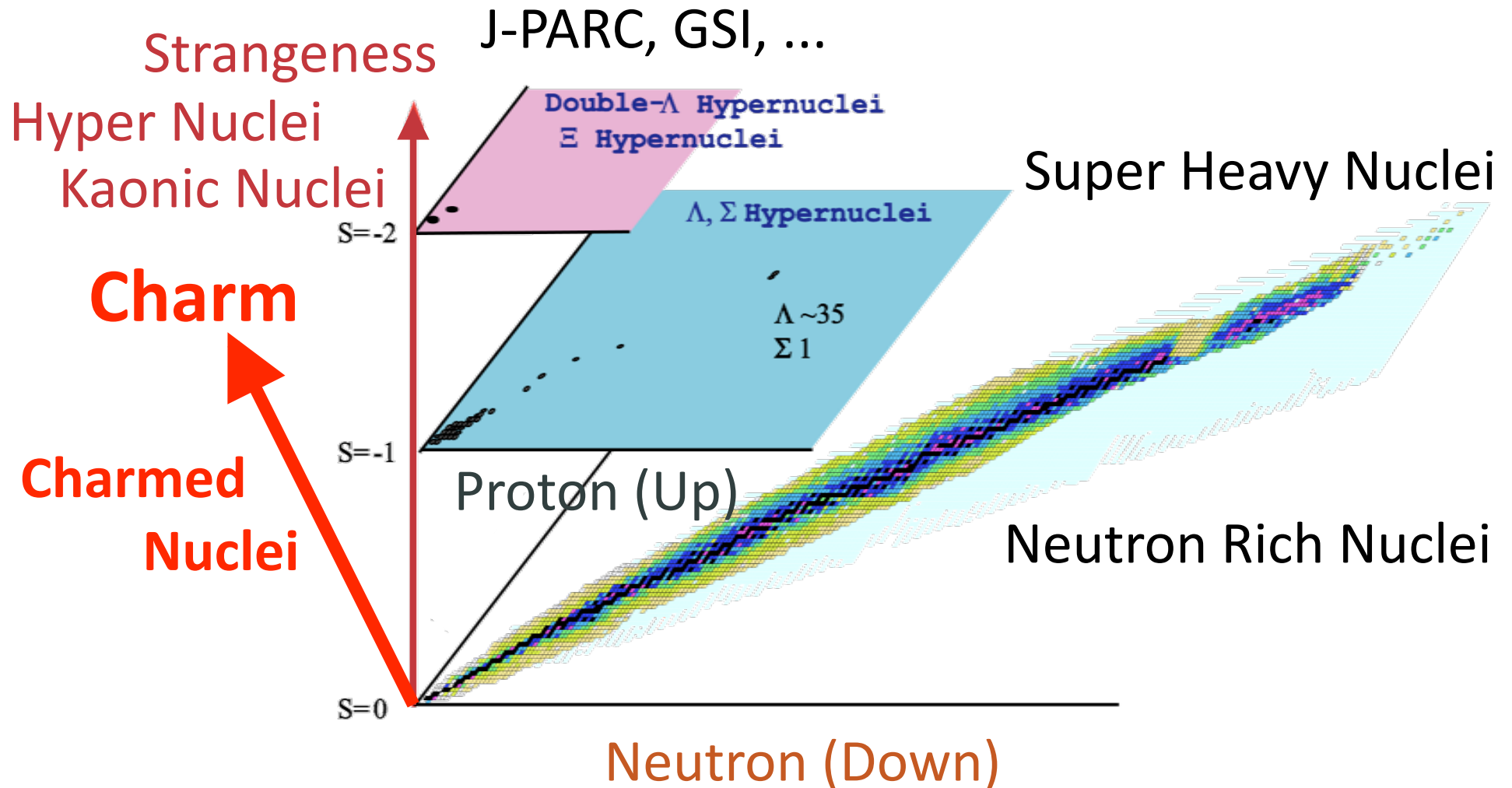


Report by PANDA@FAIR (Mar. 2009)
Krein's talk in FB19 (Sep. 2009)
Ohnishi's talk in HQP (Sep. 2009)

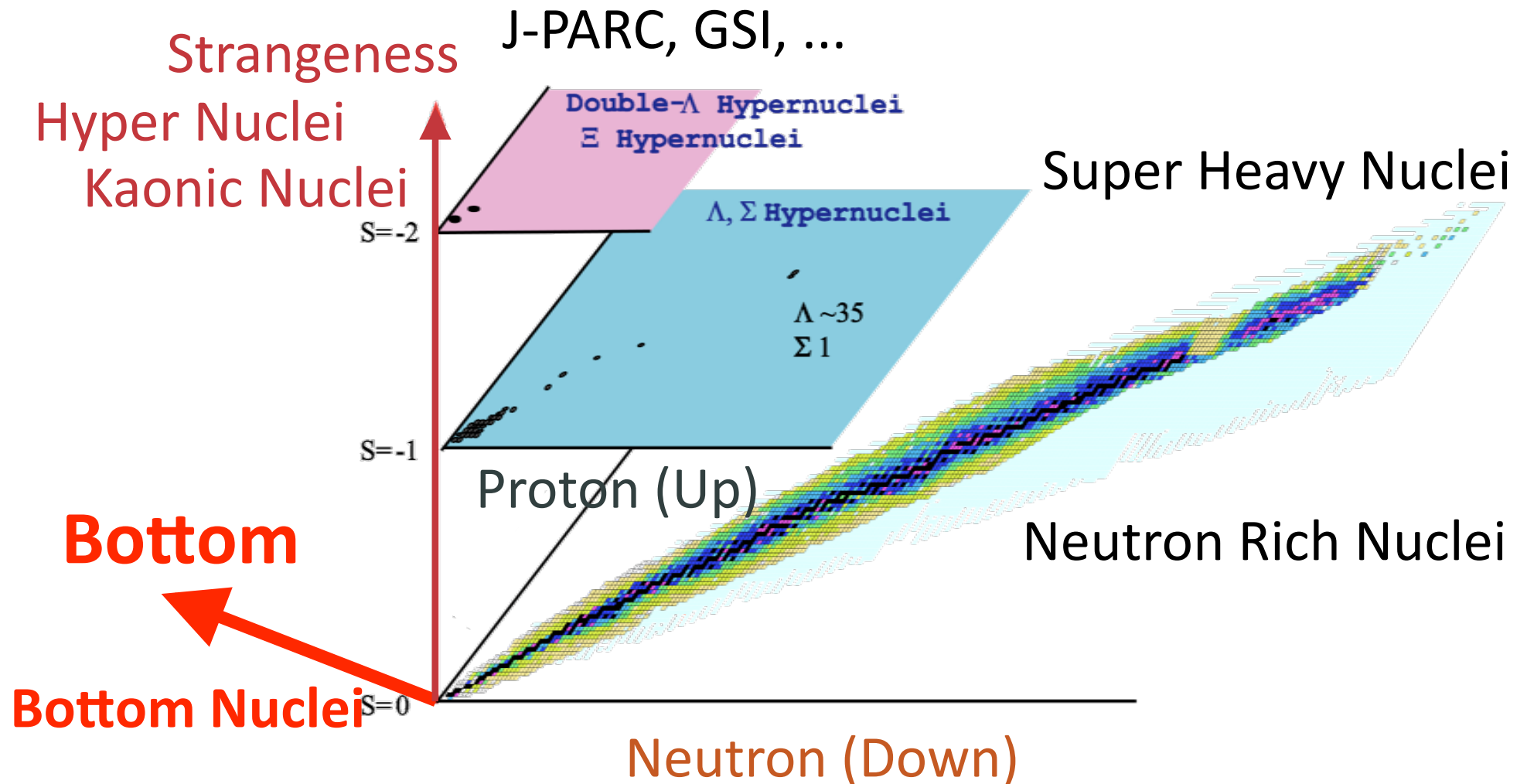
Multi-Favored Nuclear Chart



Multi-Favored Nuclear Chart

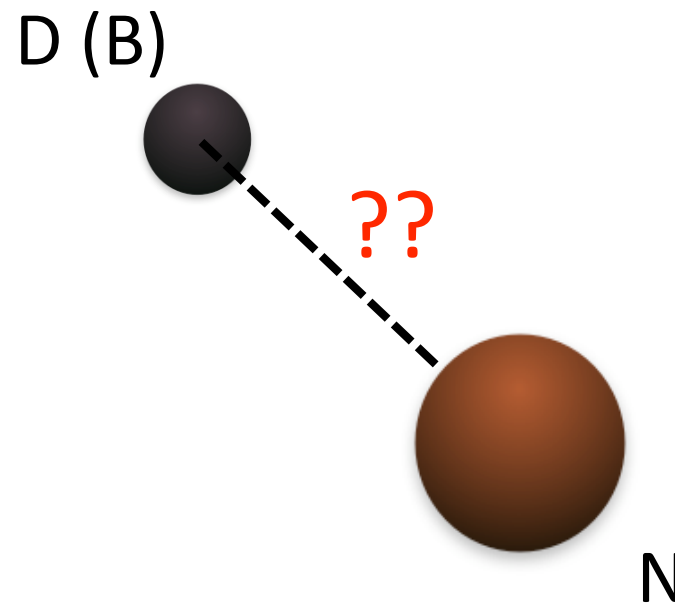


Multi-Favored Nuclear Chart



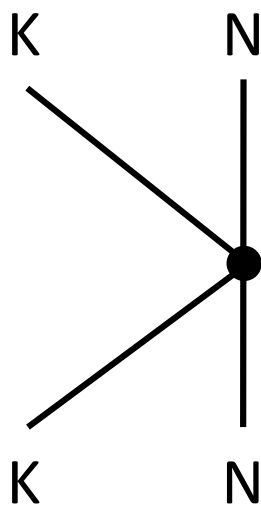
DN and BN potential

Q. What is the interaction between D (B) and N ?

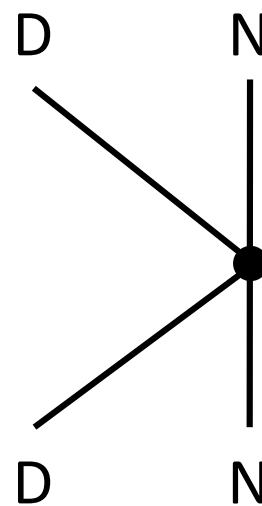


DN and BN potential

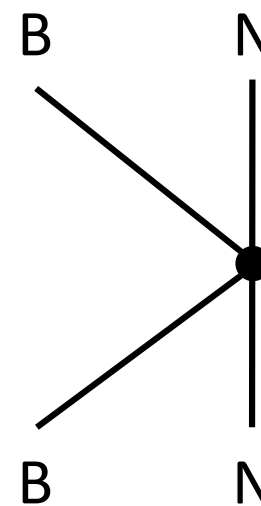
Strangeness, Charm, Bottom, ...



SU(3)



SU(4)



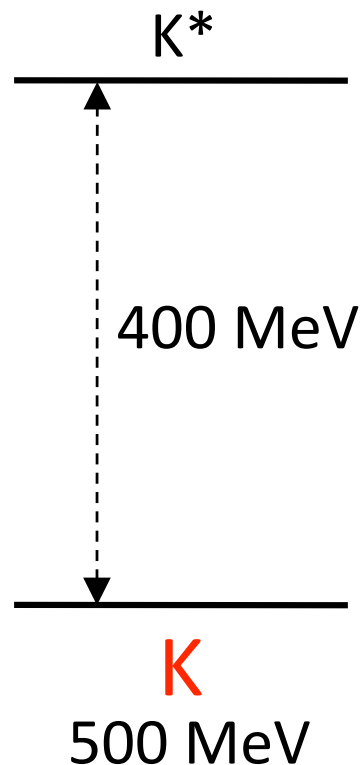
SU(5)

Weinberg-Tomozawa

Lutz-Kolomeitsev (2004), Hoffmann-Lutz (2005), Mizutani-Ramos (2006), Gamermann-Oset-Strottman-Vacas (2007), Haidenbauer-Krein-Meissner-Sibirtsev (2007), ...

DN and BN potential

Strangeness, Charm, Bottom, ...

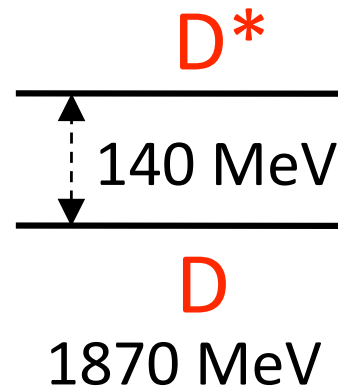


Chiral symmetry

K^* is almost irrelevant in dynamics.

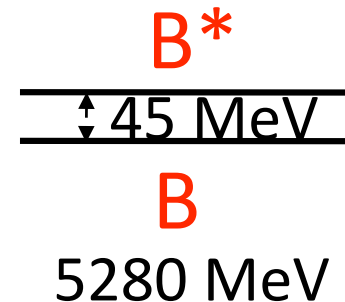


D^* and B^* are new ingredients!



D

1870 MeV



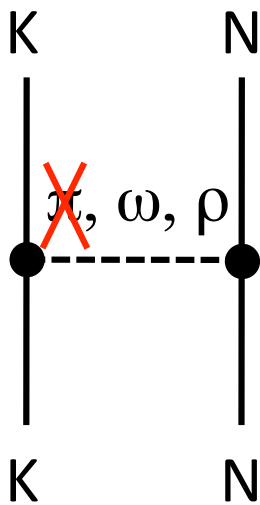
B

5280 MeV

Heavy quark symmetry

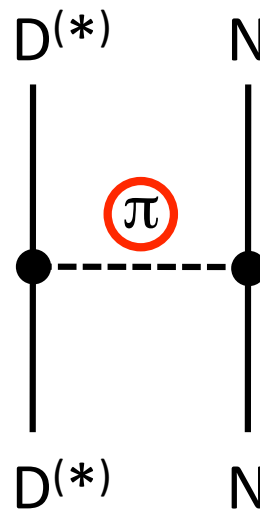
DN and BN potential

One pion exchange potential



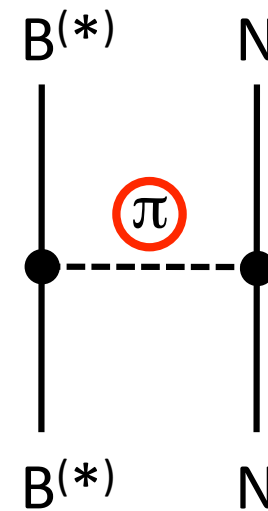
One pion is absent.
(short range force)

Weinberg-Tomozawa



One pion is present.
(long range force)

One pion exchange potential
(OPEP)

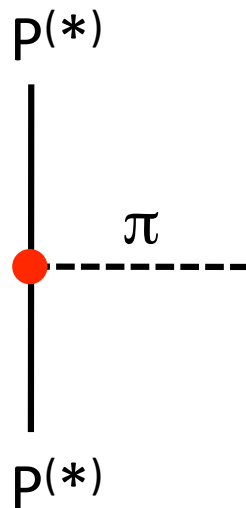


DN and BN potential

One pion exchange potential

$\rho(*)\rho(*)$ π vertex

Heavy quark symmetry



$P=D$ or B

$$\mathcal{L}_{\pi HH} = g \operatorname{tr} \bar{H}_a H_b \gamma_\nu \gamma_5 A_{ba}^\nu$$

G. Burdman and J.F. Donoghue (1992)

M.B. Wise (1992)

T.-M. Yan, H.-Y. Cheng, C.-Y. Cheung,

G.-L. Lin, Y.C. Lin and H.-L. Yu (1997)

$$\left[\begin{array}{l} H_a = \frac{1 + \not{v}}{2} [P_{a\mu}^* \gamma^\mu - P_a \gamma_5] \\ A_\mu \simeq \frac{i}{f} \partial_\mu \mathcal{M} \end{array} \right. \quad \text{with} \quad \bar{H}_a = \gamma_0 H_a^\dagger \gamma_0$$

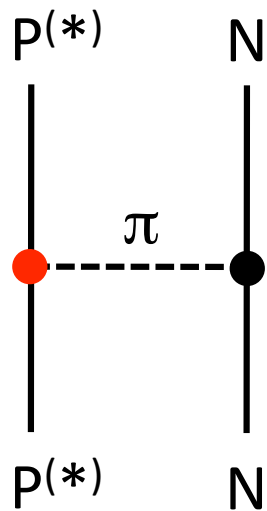
vector + pseudoscalar

$$\mathcal{M} = \begin{pmatrix} \frac{\pi^0}{\sqrt{2}} + \frac{\eta}{\sqrt{6}} & \pi^+ & K^+ \\ \pi^- & -\frac{\pi^0}{\sqrt{2}} + \frac{\eta}{\sqrt{6}} & K^0 \\ K^- & \bar{K}^0 & -\sqrt{\frac{2}{3}}\eta \end{pmatrix}$$

DN and BN potential

One pion exchange potential

S. Y. and K. Sudoh (2008)



PN-P*N and P*N-P*N potentials

$$V_{PN \rightarrow P^*N} = -\frac{gg_{\pi NN}}{\sqrt{2}m_N f} \frac{1}{4\pi} \frac{\mu^2}{3} \times \left[\vec{\varepsilon}^{(\lambda)\dagger} \cdot \vec{\sigma} C(r; \mu) + S_{\varepsilon^{(\lambda)}}^\dagger T(r; \mu) \right] \vec{\tau}_P \cdot \vec{\tau}_N$$

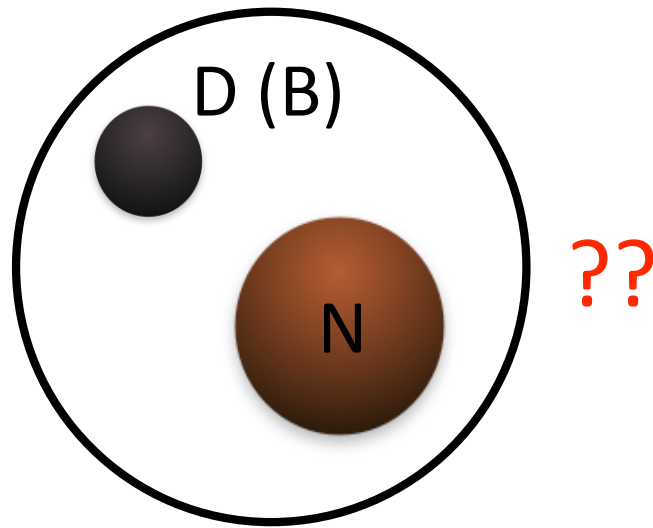
$$V_{P^*N \rightarrow P^*N} = \frac{gg_{\pi NN}}{\sqrt{2}m_N f} \frac{1}{4\pi} \frac{m_\pi^2}{3} \times \left[\vec{T} \cdot \vec{\sigma} C(r; m_\pi) + S_T T(r; m_\pi) \right] \vec{\tau}_P \cdot \vec{\tau}_N$$

$$\left[\begin{array}{l} T(r; m) = -\frac{4\pi}{m^2} \int \frac{d^3p}{(2\pi)^3} \frac{\vec{q}^2}{\vec{q}^2 + m^2} \{3(\vec{\sigma}_1 \cdot \hat{q})(\vec{\sigma}_2 \cdot \hat{q}) - \vec{\sigma}_1 \cdot \vec{\sigma}_2\} e^{i\vec{q} \cdot \vec{r}} \frac{\Lambda_N^2 - m^2}{\Lambda_N^2 + \vec{q}^2} \frac{\Lambda_P^2 - m^2}{\Lambda_P^2 + \vec{q}^2} \\ C(r; m) = \frac{4\pi}{m^2} \int \frac{d^3p}{(2\pi)^3} \frac{1}{\vec{q}^2 + m^2} e^{i\vec{q} \cdot \vec{r}} \frac{\Lambda_N^2 - m^2}{\Lambda_N^2 + \vec{q}^2} \frac{\Lambda_P^2 - m^2}{\Lambda_P^2 + \vec{q}^2} \end{array} \right.$$

DN and BN bound states

DN and BN bound states

Q. Are there bound states of D (B) and N ?



DN and BN bound states

Classification of states

S. Y. and K. Sudoh (2008)

$$|\text{state}\rangle = |PN\rangle + |P^*N\rangle \quad P = D^{\text{bar}}(c^{\text{bar}}q), \quad B(b^{\text{bar}}q)$$

* No annihilation process

DN and BN bound states

Classification of states

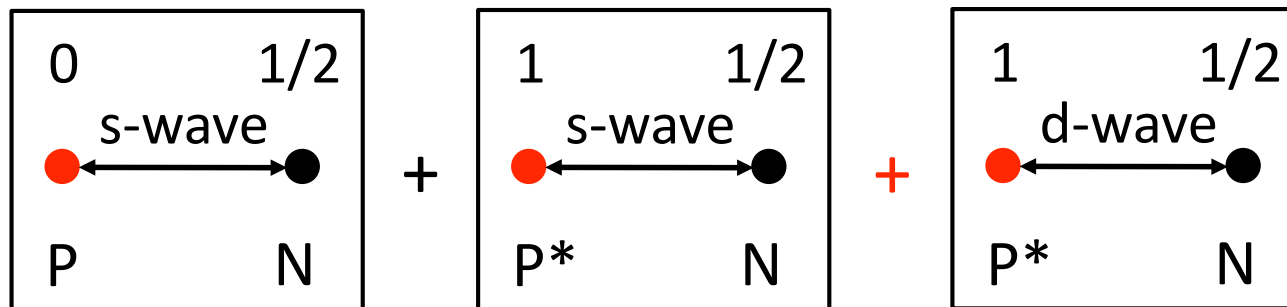
S. Y. and K. Sudoh (2008)

$$|\text{state}\rangle = |PN\rangle + |P^*N\rangle \quad P = D^{\text{bar}}(c^{\text{bar}}q), B(b^{\text{bar}}q)$$

* No annihilation process

$$\underline{J^P = 1/2^- \quad (l=0 \text{ or } 1)}$$

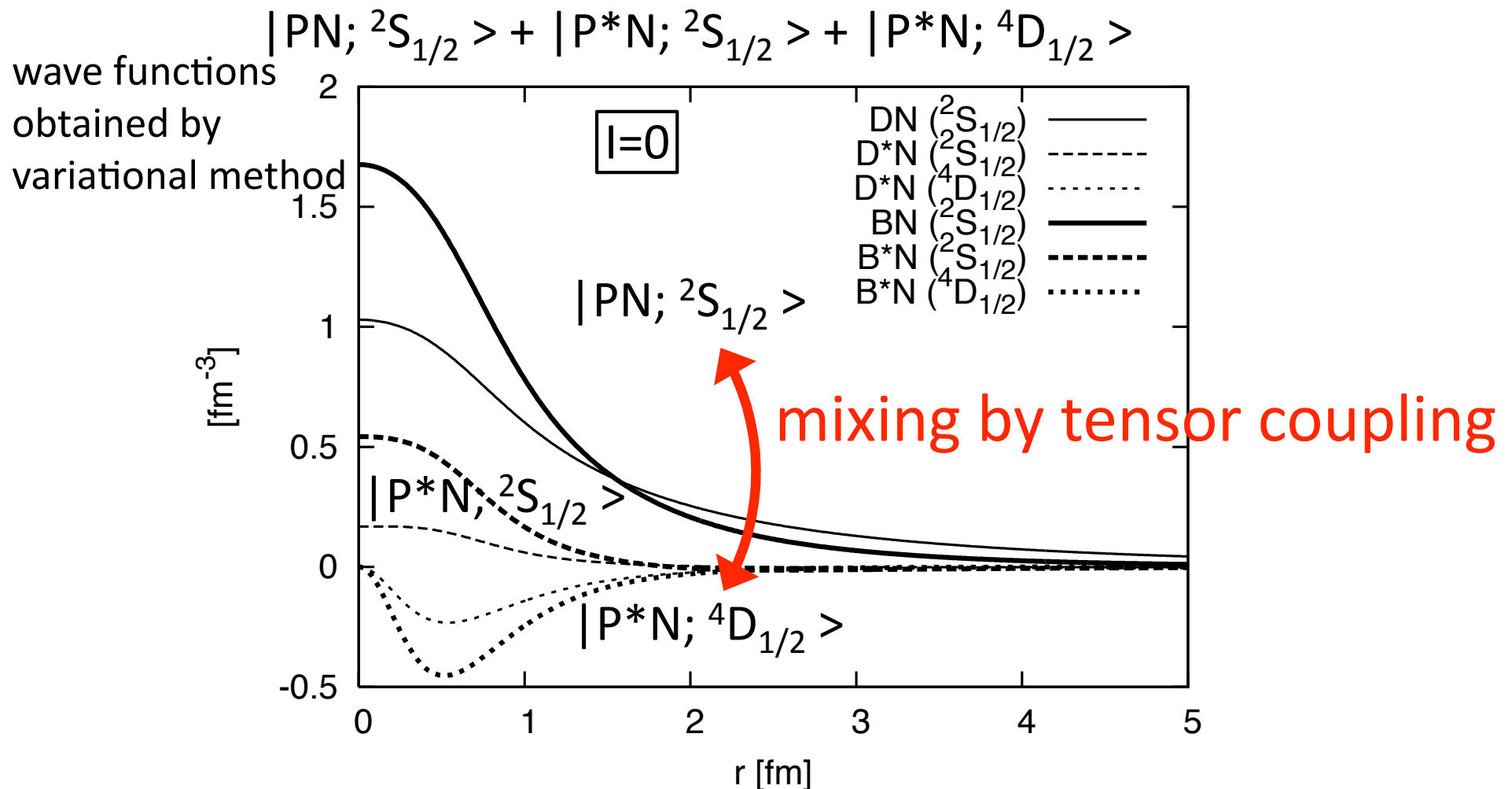
$$|PN; {}^2S_{1/2}\rangle + |P^*N; {}^2S_{1/2}\rangle + |P^*N; {}^4D_{1/2}\rangle$$



DN and BN bound states

$J^P=1/2^-$ ($l=0$ or 1)

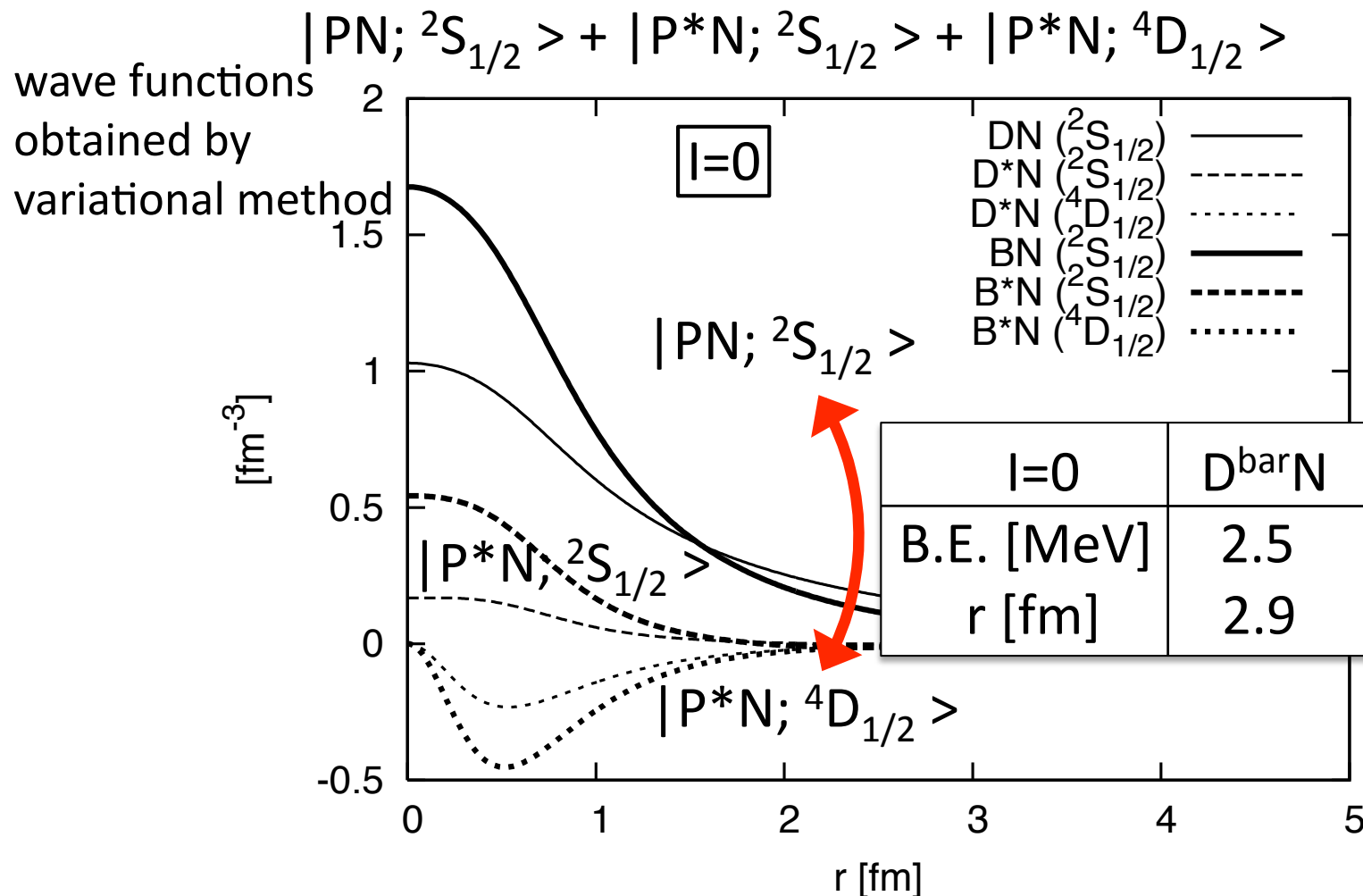
S. Y. and K. Sudoh (2008)



DN and BN bound states

$J^P=1/2^-$ ($l=0$ or 1)

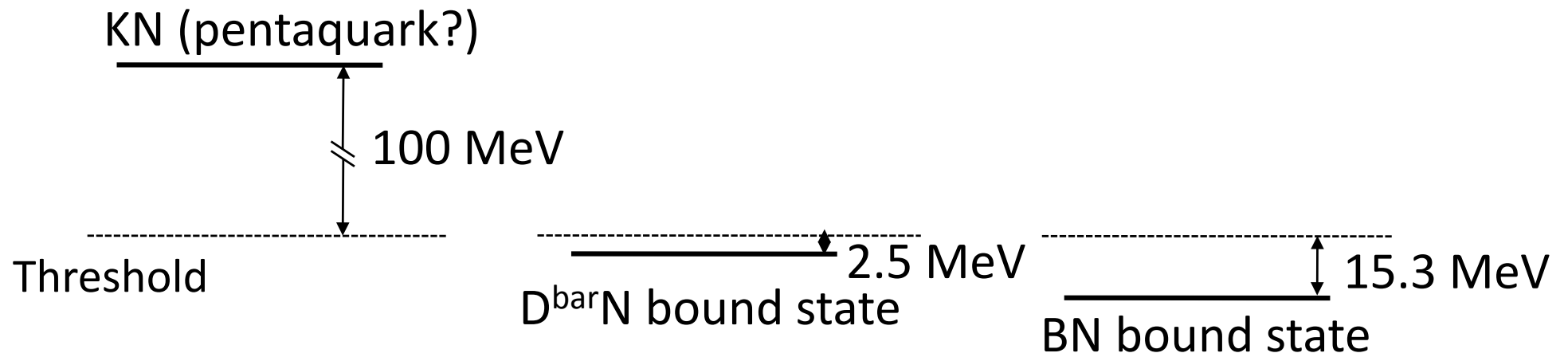
S. Y. and K. Sudoh (2008)



DN and BN bound states

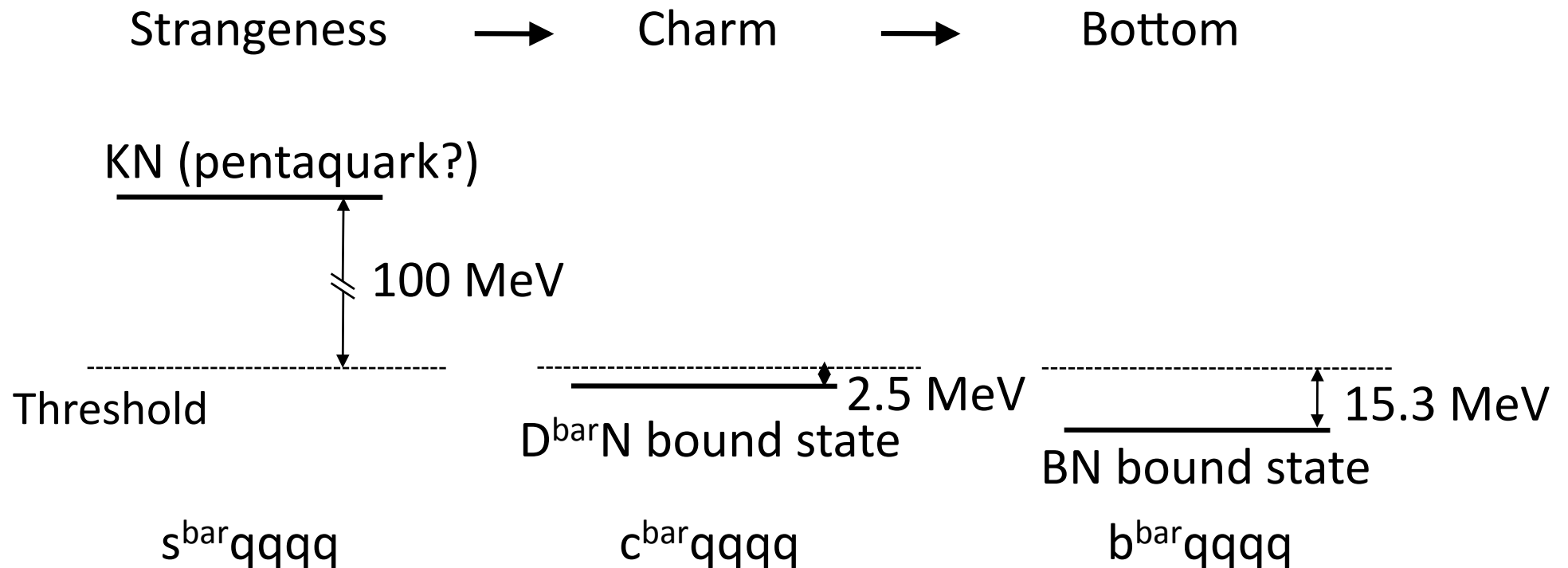
Transition from strangeness to charm and bottom systems?

Strangeness $\longrightarrow$ Charm $\longrightarrow$ Bottom



DN and BN bound states

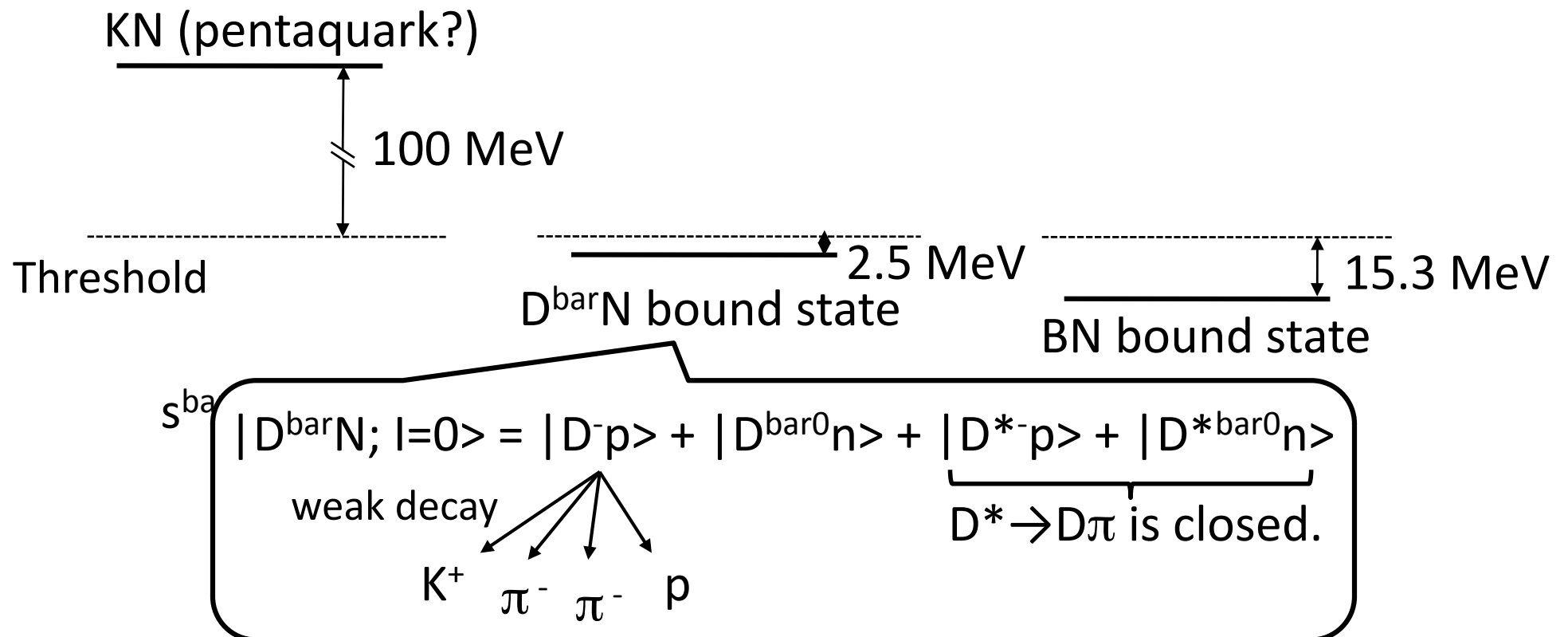
Transition from strangeness to charm and bottom systems?



DN and BN bound states

Transition from strangeness to charm and bottom systems?

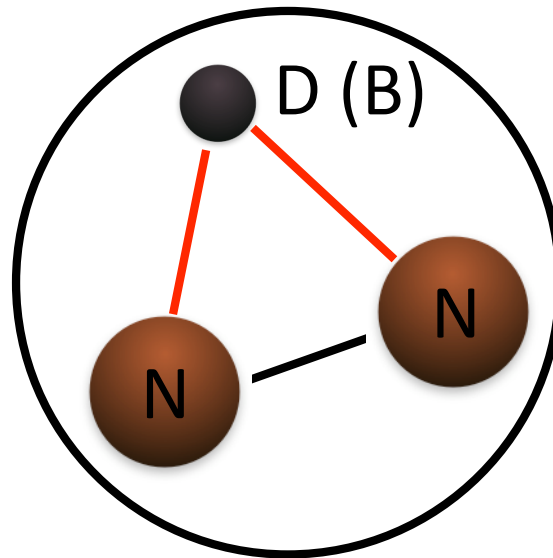
Strangeness $\rightarrow$ Charm $\rightarrow$ Bottom



Exotic charm/bottom nuclei

Exotic charm/bottom nuclei

Q. Are there exotic nuclei with D (B) ?

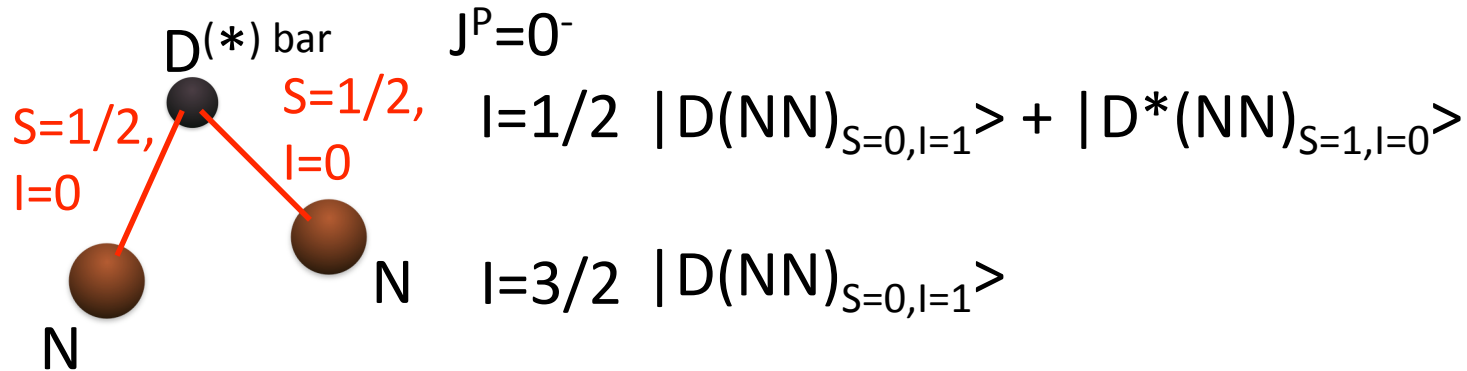


A=2 charmed (bottom) nucleus ?

Exotic charm/bottom nuclei

$D^{(*)\bar{}}NN$ and $B^{(*)}NN$ systems

Classification of states



$J^P = 1^-$

$I = 1/2 \quad |D(NN)_{S=1, I=0}\rangle + |D^*(NN)_{S=1, I=0}\rangle + |D^*(NN)_{S=0, I=1}\rangle$

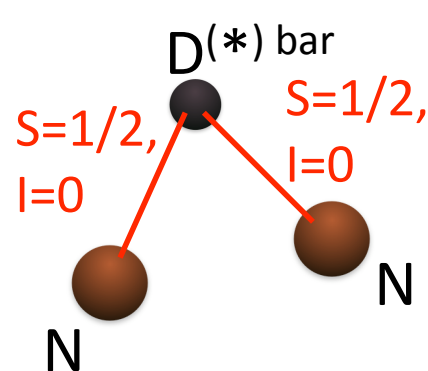
$I = 3/2 \quad |D(NN)_{S=0, I=1}\rangle$

Exotic charm/bottom nuclei

$D^{(*)\bar{}}NN$ and $B^{(*)}NN$ systems

Classification of states

Fraction of $(D^{(*)\bar{}}N)_{S=1/2, I=0}$



$J^P=0^-$

$$I=1/2 \quad |D(NN)_{S=0, I=1}\rangle + |D^*(NN)_{S=1, I=0}\rangle$$

$\frac{3}{4}$ $\frac{1}{4}$

$$I=3/2 \quad |D(NN)_{S=0, I=1}\rangle$$

0

$J^P=1^-$

$$I=1/2 \quad |D(NN)_{S=1, I=0}\rangle + |D^*(NN)_{S=1, I=0}\rangle + |D^*(NN)_{S=0, I=1}\rangle$$

$\frac{1}{4}$ $\frac{1}{6}$ $\frac{1}{4}$

$$I=3/2 \quad |D(NN)_{S=0, I=1}\rangle$$

0

Conclusion & Perspective

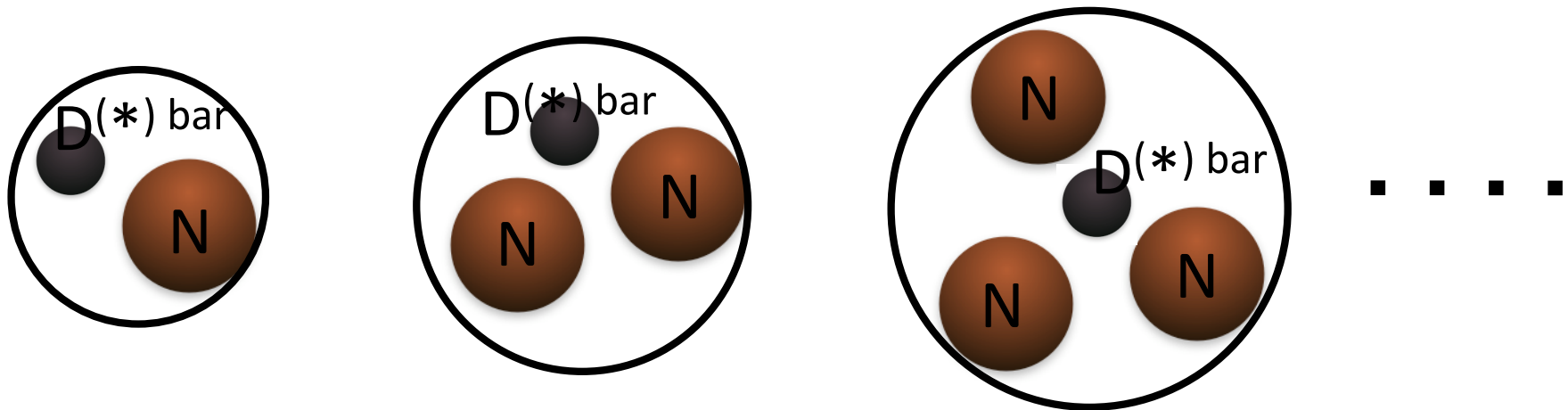
- We discuss $D^{(*)\bar{c}}N$ and $B^{(*)}N$ bound states with respect to heavy quark symmetry.

$J^P=1/2^-, l=0$	$D^{\bar{c}}N$	BN
B.E. [MeV]	2.5	15.3
r [fm]	2.9	1.4

- $DN-D^*N$ ($BN-B^*N$) mixing with tensor coupling is important.
- Charmed nuclei will be studied in experiments (J-PARC, GSI, ...).

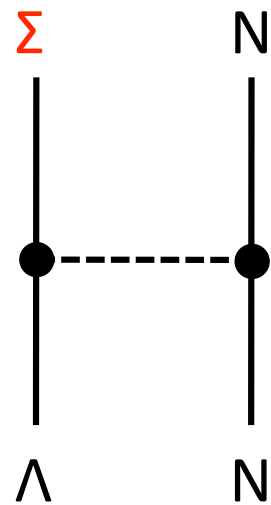
Conclusion & Perspective

- Further studies;
 - Short range potentials (under progress)
 - Phenomenological approach (vector meson/chiral loop)
 - Lattice QCD
 - $D^{(*)\text{bar}}NN$ and $B^{(*)}NN$ systems (under progress)
 - Nuclei with more baryon number (in near future)



Perspective

Hypernuclei

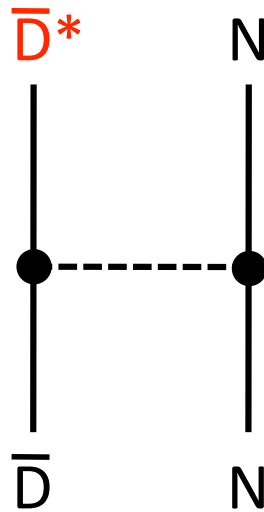


$$m_{\Lambda} \approx m_{\Sigma}$$

Λ - Σ mixing

Hiyama *et al.*

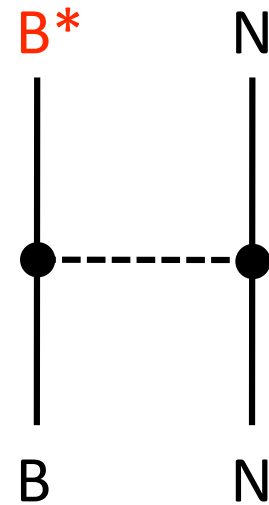
D (B) meson nuclei



$$m_D \approx m_{D^*}$$

D- D^* (B- B^*) mixing

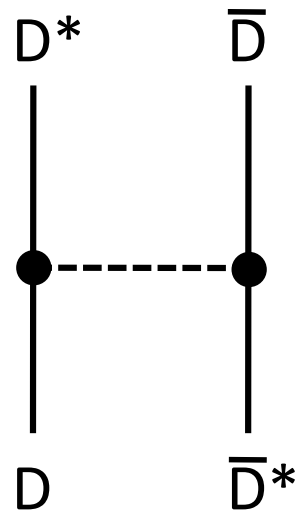
This work



$$m_B \approx m_{B^*}$$

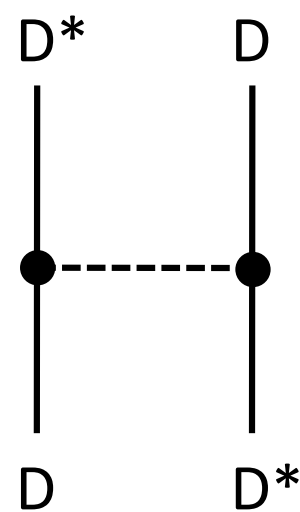
Perspective

$D-\bar{D}^*$ molecule
X(3872)



Tornqvist (1994)

$D-D^*$ molecule
 $T_{cc} (?)$



Manohar and Wise (1992)

$D-D^*$ mixing

Perspective

Diquark model; S. H. Lee and S. Y. (2009)

✓ Tetraquark T_{cc} ($c^{\text{bar}}c^{\text{bar}}ud$; $C=+2$)

T_{cc}^1	$ud\bar{c}\bar{c}$	$us\bar{c}\bar{c}$	$ds\bar{c}\bar{c}$
	-74.9 $\bar{D}^0 + D^{*-}, \bar{D}^{*0} + D^-$	-4.3 $\bar{D}^0 + D_s^{*-}$	-4.3 $D^- + D_s^{*-}$
T_{bb}^1	$ud\bar{b}\bar{b}$	$us\bar{b}\bar{b}$	$ds\bar{b}\bar{b}$
	-123.8 $B^+ + B^{*0}, B^{*+} + B^0$	-61.4 $B^+ + B_s^{*0}$	-61.4 $B^0 + B_s^{*0}$

MeV

✓ Pentaquark Θ_{cs} ($udusc^{\text{bar}}$; $S=-1, C=+1$)

	config. of Θ	config. of $M + B$	$\bar{u}$	$\bar{s}$	$\bar{c}$	$\bar{b}$
Θ_{qs}	$udus\bar{q}$	$uds+u\bar{q}$	389.4	198.9	8.4	-56.4
Θ_{qss}	$usds\bar{q}$	$ds s+u\bar{q}$	389.4	198.9	8.4	-56.4
		$uds+s\bar{q}$	256.8	142.5	28.4	-10.7

MeV

Perspective

Diquark model; S. H. Lee and S. Y. (2009)

✓ Tetraquark T_{cc} ($c^{\text{bar}}c^{\text{bar}}ud$; C=+2)

T_{cc}^1	$ud\bar{c}\bar{c}$	$us\bar{c}\bar{c}$	$ds\bar{c}\bar{c}$
	-74.9 $\bar{D}^0 + D^{*-}, \bar{D}^{*0} + D^-$	-4.3 $\bar{D}^0 + D_s^{*-}$	-4.3 $D^- + D_s^{*-}$
T_{bb}^1	$ud\bar{b}\bar{b}$	$us\bar{b}\bar{b}$	$ds\bar{b}\bar{b}$
	-123.8 $B^+ + B^{*0}, B^{*+} + B^0$	-61.4 $B^+ + B_s^{*0}$	-61.4 $B^0 + B_s^{*0}$

MeV

✓ Pentaquark Θ_{cs} ($ud\bar{u}s\bar{c}$; S=-1, C=+1)

	config. of Θ	config. of $M + B$	$\bar{u}$	$\bar{s}$	$\bar{c}$	$\bar{b}$
Θ_{qs}	$ud\bar{u}s\bar{q}$	$ud\bar{s} + u\bar{q}$	389.4	198.9	8.4	-56.4
Θ_{qss}	$us\bar{d}s\bar{q}$	$ds\bar{s} + u\bar{q}$	389.4	198.9	8.4	-56.4
		$ud\bar{s} + s\bar{q}$	256.8	142.5	28.4	-10.7

MeV

✓ H dibaryon H_c (udusuc)

Diquark model; S. H. Lee and S. Y. (2009)

H

H_c

H_b

H_{cb}

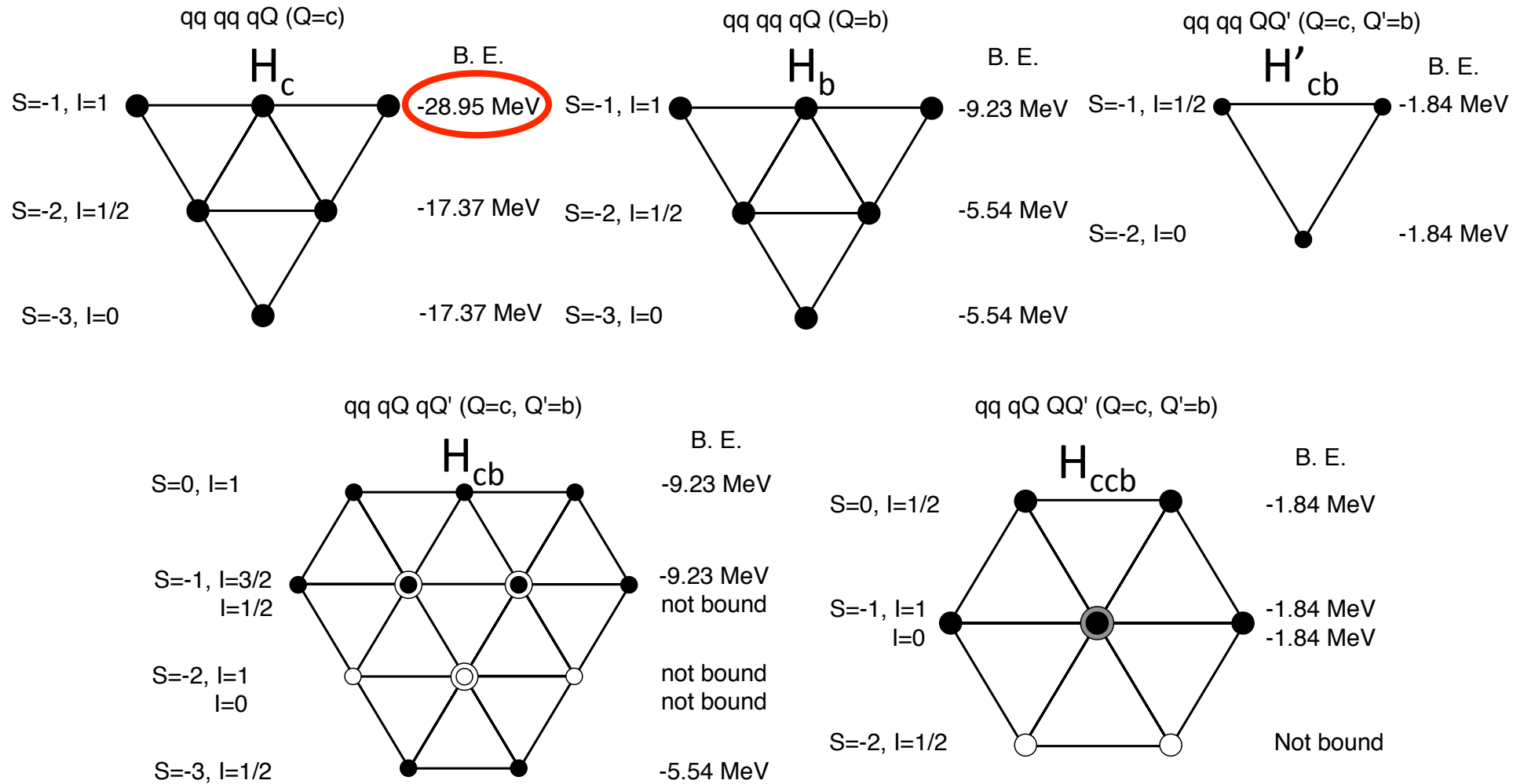
H'_{cb}

H_{ccb}

	flavor	config. of H	I	I_z	S	config. of $B + B'$	B_H [MeV]
$qq\,qq\,qq$	22200	$ud\,us\,ds$	0	0	-2	$ud\,s + ud\,s$	-28.95
$qq\,qq\,qQ$ ($Q = c$)	31110	$ud\,us\,uc$	1	1	-1	$ud\,c + us\,u, ud\,u + us\,c$	-28.95
	22110	$ud\,\frac{1}{\sqrt{2}}(us\,dc + ds\,uc)$	1	0	-1	$ud\,s + ud\,c, ud\,s + ud\,c$	
	13110	$ud\,ds\,dc$	1	-1	-1	$ud\,c + ds\,d, ud\,d + ds\,c$	-17.37
	21210	$ud\,us\,sc$	1/2	1/2	-2	$ud\,c + us\,s, ud\,s + us\,c$	
	12210	$ud\,ds\,sc$	1/2	-1/2	-2	$ud\,c + ds\,s, ud\,s + ds\,c$	
	11310	$\frac{1}{\sqrt{2}}(us\,ds - ds\,us)\,sc$	0	0	-3	$us\,s + ds\,c, us\,c + ds\,s$	-17.37
$qq\,qq\,qQ$ ($Q = b$)	31101	$ud\,us\,ub$	1	1	-1	$ud\,u + us\,b, udb + us\,u$	-9.23
	22101	$ud\,\frac{1}{\sqrt{2}}(us\,db + ds\,ub)$	1	0	-1	$ud\,s + ud\,b, ud\,s + udb$	
	13101	$ud\,ds\,db$	1	-1	-1	$udb + ds\,d, udd + ds\,b$	-5.54
	21202	$ud\,us\,sb$	1/2	1/2	-2	$udb + us\,s, ud\,s + us\,b$	
	12201	$ud\,ds\,sb$	1/2	-1/2	-2	$udb + ds\,s, ud\,s + ds\,b$	
	11301	$\frac{1}{\sqrt{2}}(us\,ds - ds\,us)\,sb$	0	0	-3	$us\,b + ds\,s, us\,s + ds\,b$	-5.54
$qq\,qQ\,qQ'$ ($Q, Q' = c, b$)	31011	$ud\,uc\,ub$	1	1	0	$udb + uc\,u, ud\,u + uc\,b$	-9.23
	22011	$ud\,\frac{1}{\sqrt{2}}(uc\,db + dc\,ub)$	1	0	0	$udb + uc\,u, ud\,u + uc\,b$	
	13011	$ud\,dc\,db$	1	-1	0	$udb + dc\,d, udd + dc\,b$	-9.23
	30111	$us\,uc\,ub$	3/2	3/2	-1	$us\,u + uc\,b, us\,b + uc\,u$	
	21111	$\frac{1}{\sqrt{3}}((us\,dc + ds\,uc)\,ub + us\,uc\,db)$	3/2	1/2	-1	$udb + us\,c, udc + us\,b$	
	12111	$\frac{1}{\sqrt{3}}((us\,dc + ds\,uc)\,db + ds\,dc\,ub)$	3/2	-1/2	-1	$udb + ds\,c, udc + ds\,b$	
	03111	$ds\,dc\,db$	3/2	-3/2	-1	$ds\,b + dc\,d, ds\,d + dc\,b$	-5.54
	10311	$us\,sc\,sb$	1/2	1/2	-3	$us\,b + sc\,s, us\,s + sc\,b$	
	01311	$ds\,sc\,sb$	1/2	-1/2	-3	$ds\,b + sc\,s, ds\,s + sc\,b$	
$qq\,qq\,QQ'$ ($Q, Q' = c, b$)	21111	$ud\,us\,cb$	1/2	1/2	-1	$udb + us\,c, udc + us\,b$	-1.84
	12111	$ud\,ds\,cb$	1/2	-1/2	-1	$udb + ds\,c, udc + ds\,b$	-1.84
	11211	$\frac{1}{\sqrt{2}}(us\,ds - ds\,us)\,cb$	0	0	-2	$us\,b + ds\,c, us\,c + ds\,b$	
$qq\,qQ\,QQ'$ ($Q, Q' = c, b$)	21021	$ud\,uc\,cb$	1/2	1/2	0	$udb + uc\,c, udc + uc\,b$	-1.84
	12021	$ud\,dc\,cb$	1/2	-1/2	0	$udb + dc\,c, udc + dc\,b$	-1.84
	20121	$us\,uc\,cb$	1	1	-1	$us\,b + uc\,c, us\,c + uc\,b$	
	11121	$\frac{1}{\sqrt{2}}(us\,dc + ds\,uc)\,cb$	1	0	-1	$us\,b + dc\,c, us\,c + dc\,b$	
	02121	$ds\,dc\,cb$	1	-1	-1	$ds\,b + dc\,c, ds\,c + dc\,b$	
	11121	$ud\,sc\,cb$	0	0	-1	$udb + sc\,c, udc + sc\,b$	

MeV

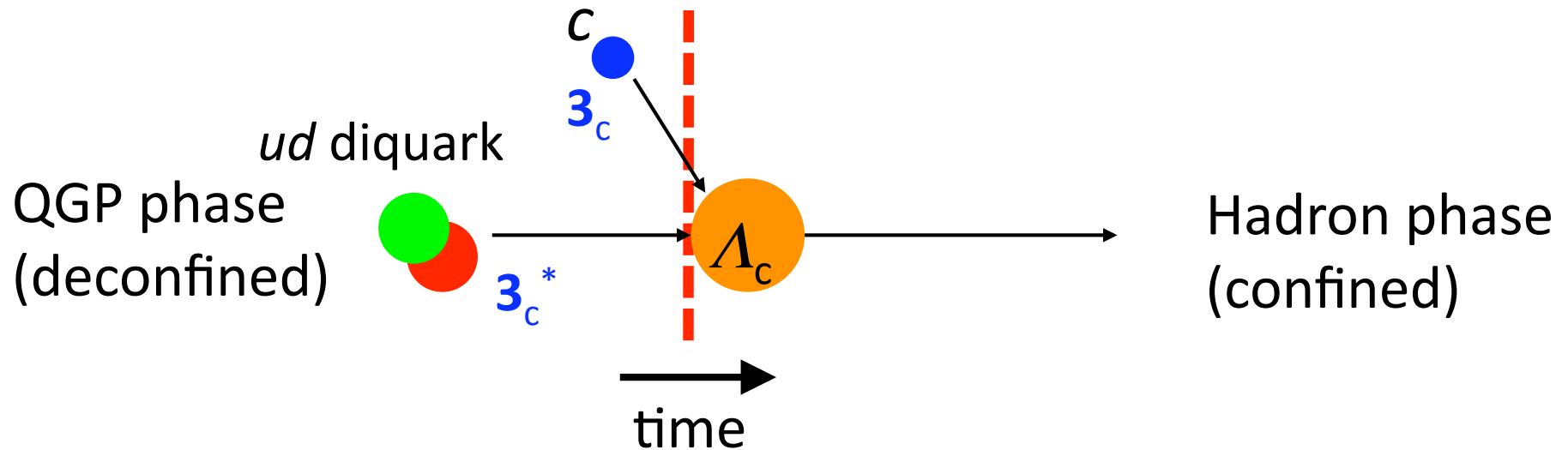
✓ H dibaryon H_c (udusuc)
multiplets of $SU(3)_f$.



→ H dibaryons ($H, H_c, H_b, \dots$) are stable as 3_f multiplet of $SU(3)_f$.

Perspective

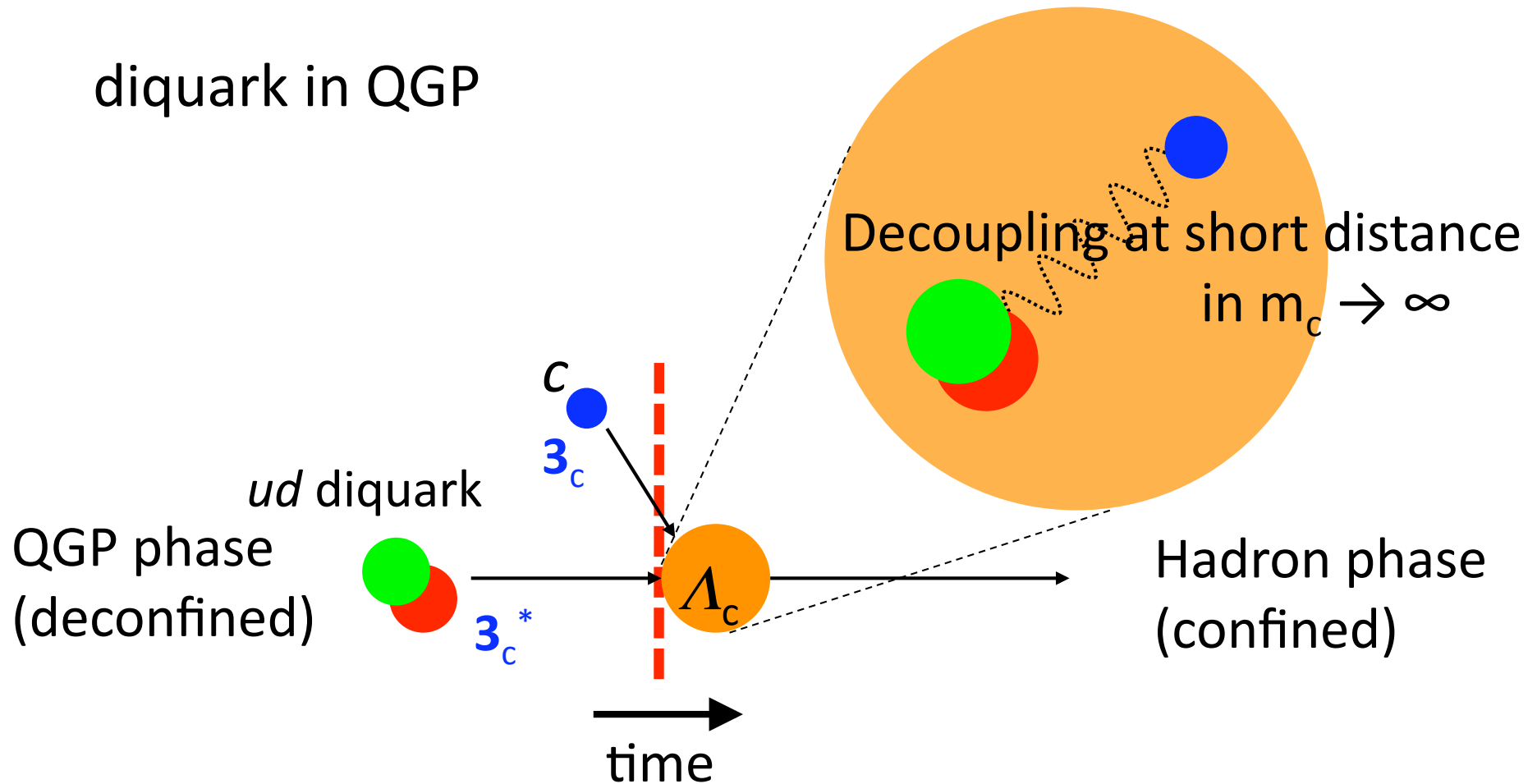
diquark in QGP



S. H. Lee, K. Ohnishi, S. Y., I. K. Yoo, C. M. Ko (2008)

S. H. Lee, Y. Oh, C. M. Ko, S. Y. (2009)

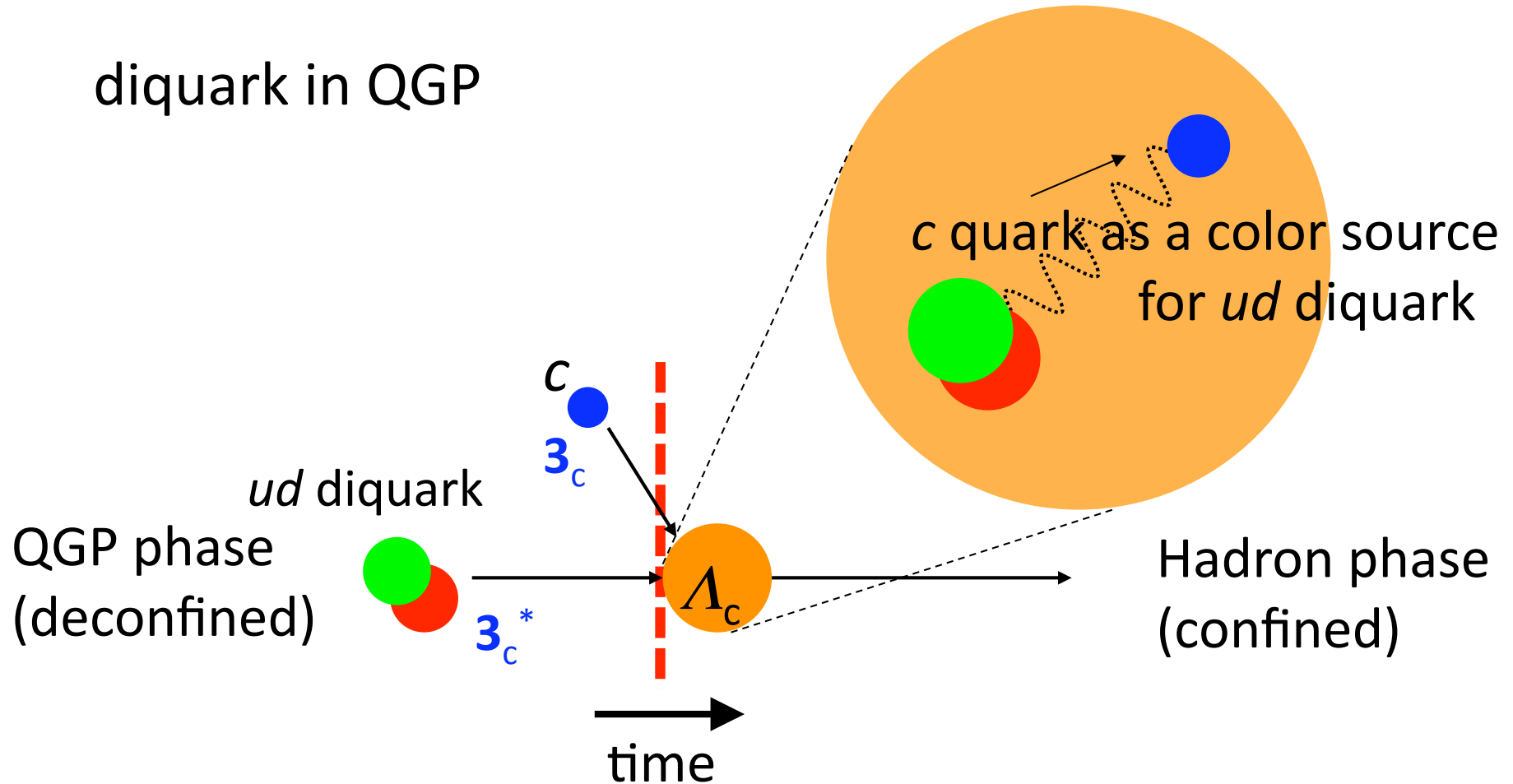
Perspective



S. H. Lee, K. Ohnishi, S. Y., I. K. Yoo, C. M. Ko (2008)

S. H. Lee, Y. Oh, C. M. Ko, S. Y. (2009)

Perspective

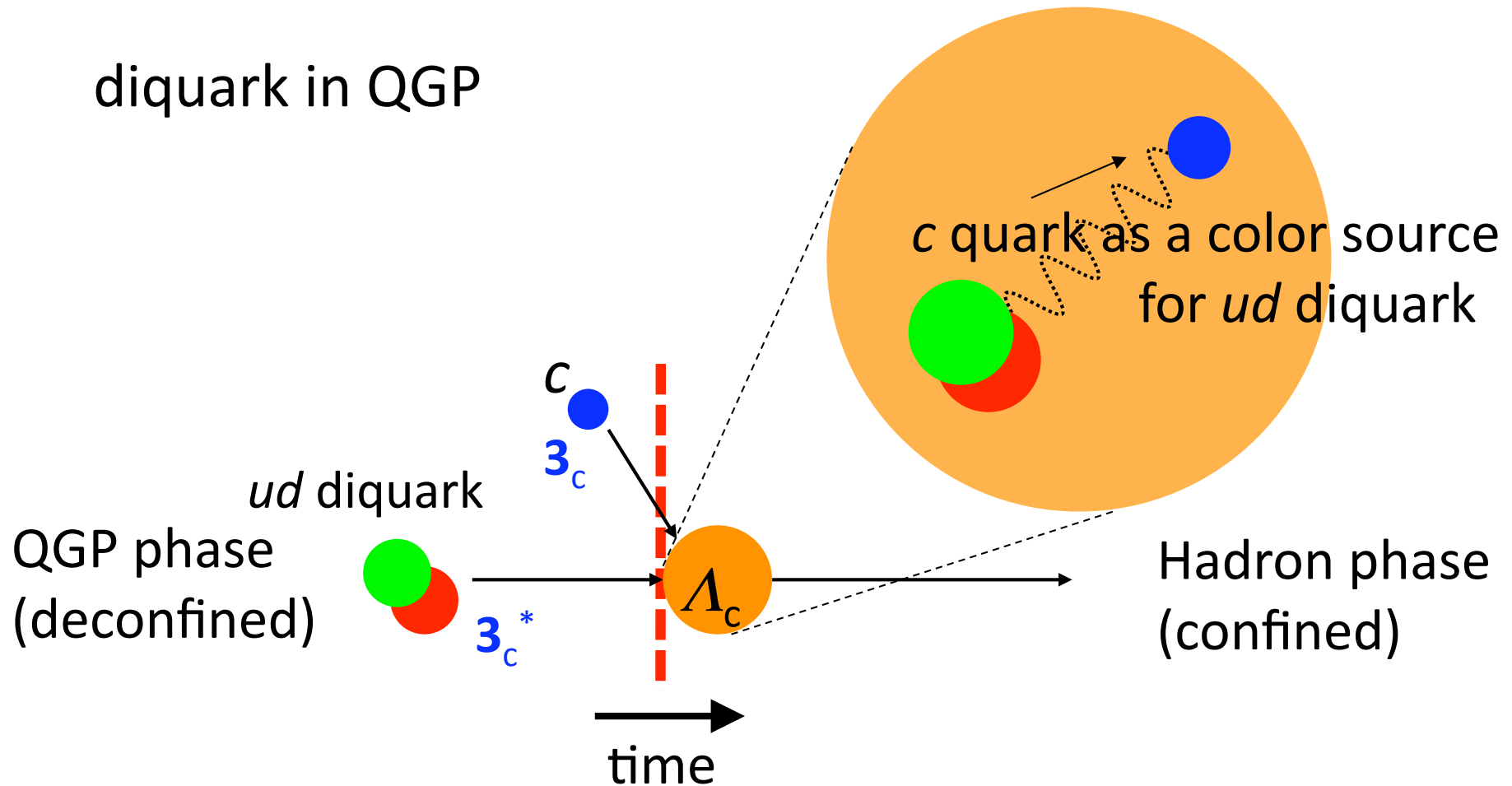


S. H. Lee, K. Ohnishi, S. Y., I. K. Yoo, C. M. Ko (2008)

S. H. Lee, Y. Oh, C. M. Ko, S. Y. (2009)

Perspective

diquark in QGP



✓ *ud* diquark $\longleftrightarrow \Lambda_c$

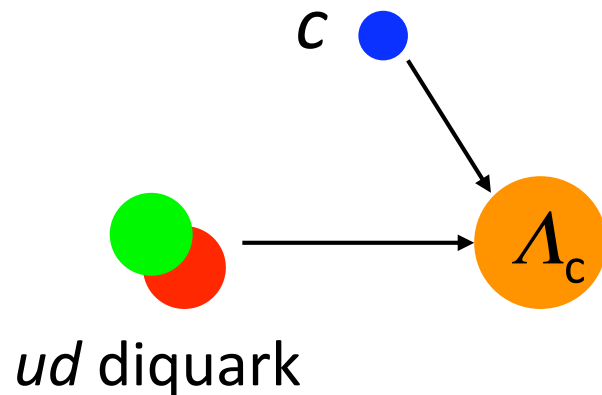
S. H. Lee, K. Ohnishi, S. Y., I. K. Yoo, C. M. Ko (2008)

S. H. Lee, Y. Oh, C. M. Ko, S. Y. (2009)

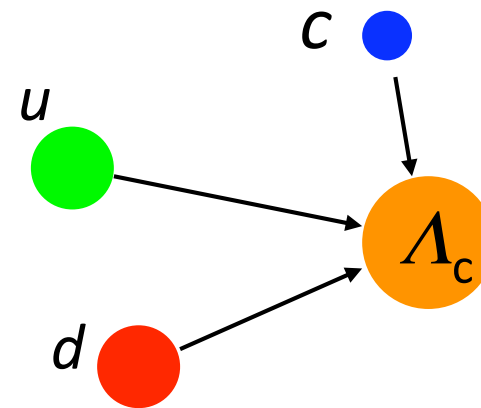
Perspective

In the process of yielding Λ_c ...

a) With diquark



b) Without diquark



Kinematically

2 body collision

>>

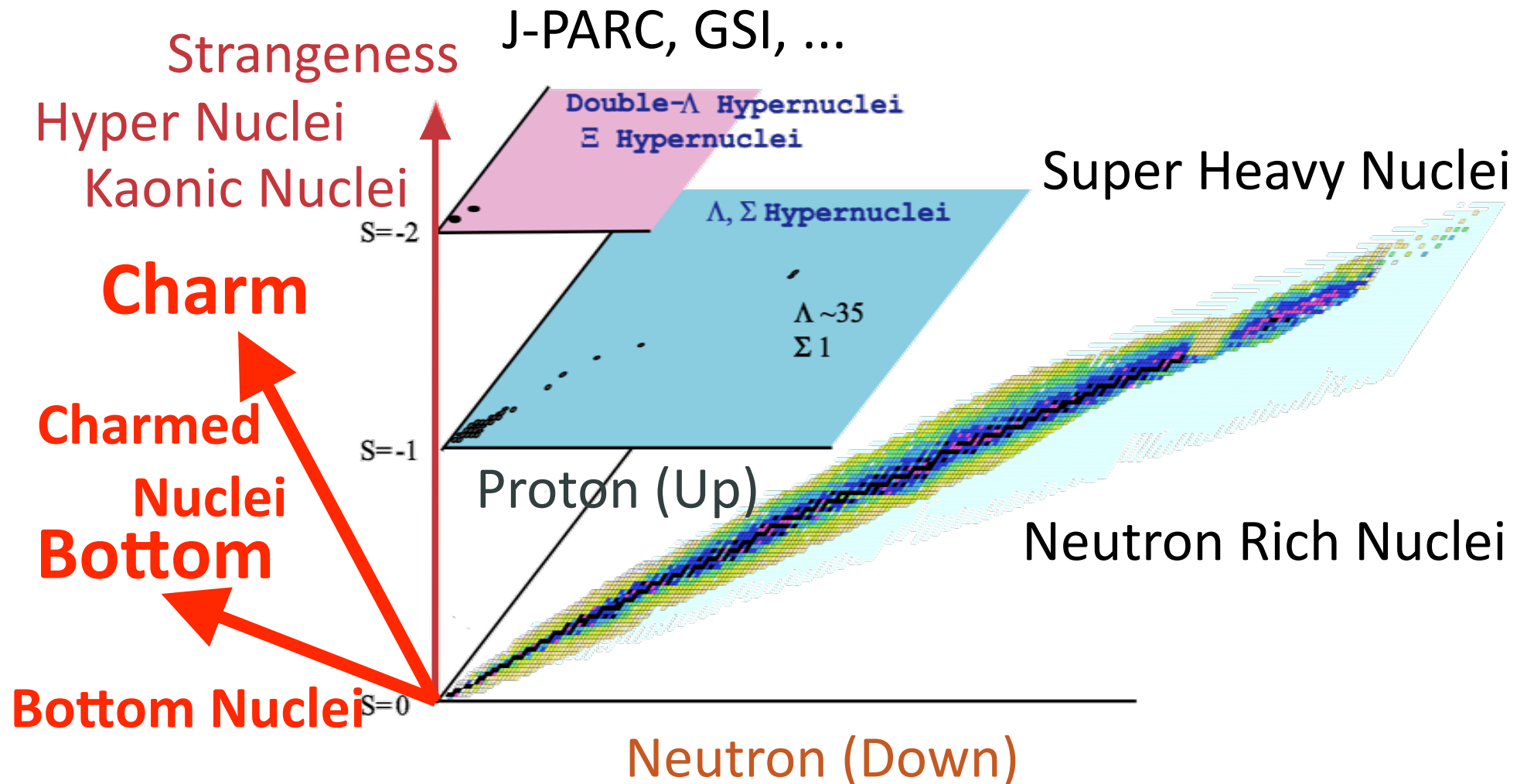
3 body collision

→ ud diquark enhances Λ_c yield by more than 2 times.

S. H. Lee, K. Ohnishi, S. Y., I. K. Yoo, C. M. Ko (2008)

S. H. Lee, Y. Oh, C. M. Ko, S. Y. (2009)

Multi-Favored Nuclear Chart



DN and BN bound states

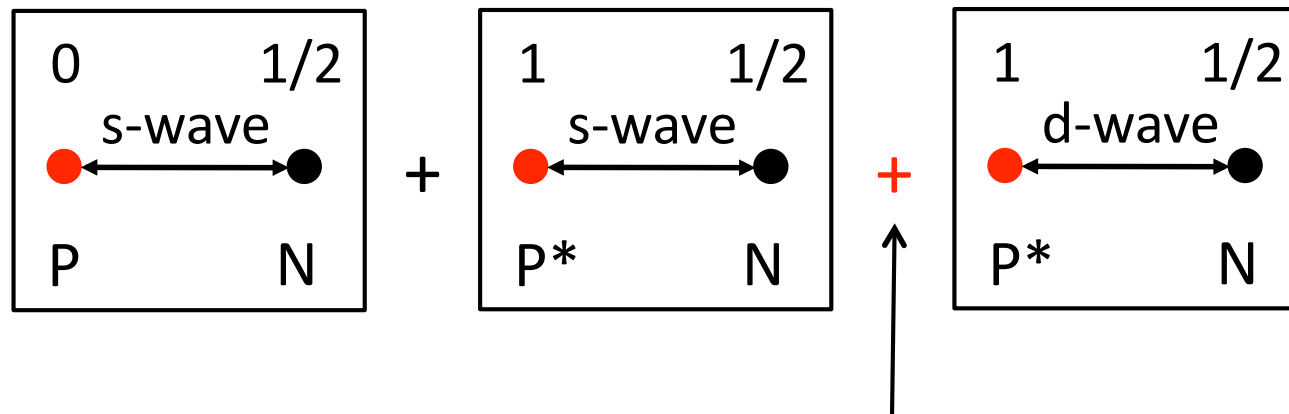
Classification of states

$$|\text{state}\rangle = |PN\rangle + |P^*N\rangle \quad P = D^{\text{bar}}(c^{\text{bar}}q), B(b^{\text{bar}}q)$$

* No annihilation process

$$\underline{J^P = 1/2^- \quad (l=0 \text{ or } 1)}$$

$$|PN; {}^2S_{1/2}\rangle + |P^*N; {}^2S_{1/2}\rangle + |P^*N; {}^4D_{1/2}\rangle$$



This was missing in Cohen-Hohler-Lebed (2005).

DN and BN bound states

Classification of states

$$|\text{state}\rangle = |\text{PN}\rangle + |\text{P}^*\text{N}\rangle \quad \text{P} = \text{D}^{\text{bar}} (c^{\text{bar}} q), \text{B} (b^{\text{bar}} q)$$

* No annihilation process

Case 1. $J^P = 1/2^-$ ($l=0$ or 1)

$$|\text{PN}; {}^2S_{1/2}\rangle + |\text{P}^*\text{N}; {}^2S_{1/2}\rangle + |\text{P}^*\text{N}; {}^4D_{1/2}\rangle$$

Case 2. $J^P = 3/2^-$ ($l=0$ or 1)

$$|\text{PN}; {}^2D_{3/2}\rangle + |\text{P}^*\text{N}; {}^4S_{3/2}\rangle + |\text{P}^*\text{N}; {}^4D_{3/2}\rangle + |\text{P}^*\text{N}; {}^2D_{3/2}\rangle$$